

BOSE EINSTEIN CONDENSATION ON INHOMOGENEOUS AMENABLE GRAPHS

FRANCESCO FIDALEO, DANIELE GUIDO, TOMMASO ISOLA

ABSTRACT. We investigate the Bose–Einstein Condensation on nonhomogeneous amenable networks for the model describing arrays of Josephson junctions. The resulting topological model, whose Hamiltonian is the pure hopping one given by the opposite of the adjacency operator, has also a mathematical interest in itself. We show that for the nonhomogeneous networks like the comb graphs, particles condensate in momentum and configuration space as well. In this case different properties of the network, of geometric and probabilistic nature, such as the volume growth, the shape of the ground state, and the transience, all play a rôle in the condensation phenomena. The situation is quite different for homogeneous networks where just one of these parameters, e.g. the volume growth, is enough to determine the appearance of the condensation.

1. INTRODUCTION

This paper is devoted to the analysis of thermodynamical states on complex networks with pure hopping Hamiltonian, in particular of those exhibiting Bose–Einstein condensation (BEC for short). Here the network is described by an infinite topological graph, and the pure hopping Hamiltonian is given, up to an additive constant, by the non diagonal part of the combinatorial Laplace operator on the graph (the hopping part).

The general setting considered here is that of Bosons on a lattice, described by Canonical Commutation Relations on a suitable dense subspace of $\ell^2(VX)$, see [2, 3] for references.

The study of BEC on infinite graphs with pure hopping Hamiltonian, in particular for the so called comb graph (see fig. 1), started with a series of papers (cf. [4] and references therein) motivated by the relevance of the comb graphs in describing physical phenomena such as arrays of Josephson junctions, and then continued in [1, 10].

As observed in [4], the simultaneous choice of an inhomogeneous graph (the comb graph) and of the pure hopping Hamiltonian can produce the so called “hidden spectrum”, namely the emergence, in the infinite volume limit, of an interval at the bottom of the spectrum of the Hamiltonian, on which the corresponding spectral measure vanishes.

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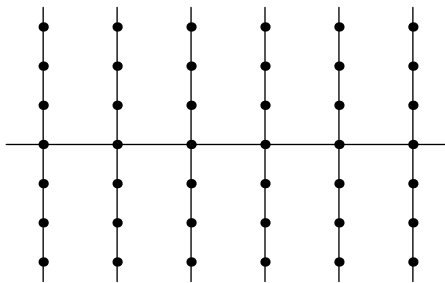


FIGURE 1. Comb graph.

Since the critical density may be seen as an integral, w.r.t. this measure, of a function with a divergence at the bottom of the spectrum, the presence of hidden spectrum immediately implies the finiteness of the critical density, even at low dimensions.

Indeed, one of the main observations of this paper concerns the decoupling, in our setting, of a series of features which equivalently characterize BEC above the critical density for homogeneous graphs with free Hamiltonian.

For periodic graphs indeed, the appearance of BEC is determined by the finiteness of the critical density. If this happens, in order to obtain a thermodynamical state with non-trivial condensate, one should choose the chemical potential μ_{Λ_n} for the cut-off region Λ_n so to keep the density constant and above the critical density.

For inhomogeneous graphs with pure hopping Hamiltonian the finiteness of the critical density is neither necessary nor sufficient for the existence of a thermodynamic limit with BEC.

Comb graphs for example have finite critical density also at low dimensions, however this does not imply the existence of locally normal thermodynamical states exhibiting BEC. This is because such finiteness is due to the presence of hidden spectrum, instead of to the integrability of the divergence at the bottom of the spectrum.

On the other hand, one can prove (cf. [6]) that for the graph \mathbb{N} a thermodynamical state with BEC exists, despite of the infinite critical density.

Moreover, higher-dimensional combs admit BEC, but the threshold dimension is not given by the growth of the volume, but by the growth of the Perron-Frobenius vector. Such threshold dimension also plays a rôle in the choice of the sequence $\{\mu_{\Lambda_n}\}$ of the finite volume chemical potentials which gives rise to locally normal states with BEC.

Apparently, the notion which is still capable of determining the existence of BEC is that of transience of the Hamiltonian operator.

The class of graphs that we analyze in the paper consists of zero density perturbations of Γ -periodic graphs with finite quotients, Γ being a discrete amenable group. We equip any such graph $X = (V, E)$ with a C^* -algebra of operators on $\ell^2(VX)$, containing in particular Γ -invariant operators with finite propagation, and endow the algebra with a finite trace τ . This trace, composed with the spectral projections of the Hamiltonian, produces the spectral measure at the infinite volume limit. Therefore the presence of hidden spectrum may be seen as a consequence of the non faithfulness of the GNS representation associated with the trace τ .

A relevant notion in this paper is that of generalized Perron–Frobenius eigenvector, namely a (not necessarily ℓ^2) vector with positive entries on VX which is an eigenvector of the adjacency operator A , with eigenvalue $\|A\|$. For the pure hopping Hamiltonian indeed, this vector describes the minimal energy, hence the Bose–Einstein condensate.

Another main technical issue studied in this paper is how small perturbations of periodic graphs modify the behavior under the thermodynamical limit. As a main tool, we prove what we call the secular equation, which relates the resolvent of the adjacency operator of a perturbed graph with that of the unperturbed one.

As a first byproduct we may prove the existence of hidden spectrum, and hence of finite critical density, for a large class of examples. Finally, we can completely analyze the thermodynamical limit for the comb graph $\mathbb{Z}^d \dashv \mathbb{Z}$, showing that locally normal KMS states at the critical density may appear exactly when $d \geq 3$.

In particular, we show that for the low dimensional combs, it is impossible to exhibit infinite volume states describing a portion of Bose–Einstein condensate having a correct local meaning. This happens since the adjacency matrix is recurrent, and we can exhibit a unique KMS state for a given inverse temperature β , which is non normal w.r.t. the Fock state. In addition, the thermodynamic limit with constant density does not describe Bose–Einstein condensate even in the transient case, i.e. for the comb graphs $\mathbb{Z}^d \dashv \mathbb{Z}$, $d \geq 3$. Such a condensate exists only at the critical density, and is obtained with a carefully chosen asymptotics of the chemical potentials for the finite volume approximations.

To end the present introduction, we recall that the comb graphs considered here are a particular case of the graphs obtained via the comb product, and that, as studied in [1, 8] and references therein, there is a correspondence between the notion of classical, resp. monotone, resp. Boolean, resp. free independence, and the notion of tensor, resp. comb, resp. star, resp. free, product of graphs.

2. GEOMETRICAL PRELIMINARIES

A *simple graph* $X = (VX, EX)$ is a collection VX of objects, called *vertices*, and a collection EX of unordered pairs of distinct vertices, called *edges*. The edge $e = \{u, v\}$ is said to join the vertices u, v , while u and v are said to be *adjacent*, which is denoted $u \sim v$.

Let us denote by $A = [A(v, w)]$, $v, w \in VX$, the adjacency matrix of X , that is,

$$A(v, w) = \begin{cases} 1 & \{v, w\} \in EX \\ 0 & \text{otherwise,} \end{cases}$$

and observe that, given VX , assigning EX is equivalent to assigning A , that is the geometrical properties of X can be expressed in terms of A . For example, a graph is connected, namely any two vertices are joined by a path, is equivalent to the irreducibility of the matrix A , the *degree* $\deg(v)$ of a vertex v , namely the number of vertices adjacent to v , is equal to (v, A^*Av) and, setting $d := \sup_{v \in VX} \deg(v)$, we have $\sqrt{d} \leq \|A\| \leq d$, namely A is bounded if and only if X has bounded degree. We denote by $D = [D(v, w)]$ the degree matrix of X , that is,

$$D(v, w) = \begin{cases} \deg v & v = w \\ 0 & \text{otherwise.} \end{cases}$$

The Laplacian on the graph is $\Delta = D - A$, so that $(\Delta f)(v) = \sum_{w \sim v} (f(w) - f(v))$, for any $f \in \ell^2(VX)$, $v \in VX$.¹

Assume now the simple graph X to be countable and with bounded degree. In the present paper we only deal with bounded operators acting on $\ell^2(VX)$, if it is not otherwise specified.

Definition 2.1. Let X be a countably infinite graph. An increasing exhaustion $\{K_n : n \in \mathbb{N}\}$ of finite subgraphs of X is called an *amenable exhaustion* of X if, setting $\mathcal{F}K_n := \{v \in VK_n : d(v, VX \setminus VK_n) = 1\}$, then $\lim_{n \rightarrow \infty} \frac{|\mathcal{F}K_n|}{|VK_n|} = 0$.

X is called an *amenable graph* if it possesses an amenable exhaustion.

We say that an operator A acting on $\ell^2(VX)$ has *finite propagation* if there exists a constant $r = r(A) > 0$ such that, for any $v \in X$, the support of Av is contained in the (closed) ball $B(v, r)$ centered in x and with radius r . It is not difficult to show that finite propagation operators form a $*$ -algebra, and we denote by $\mathcal{A}_{FP}(X)$ the generated C^* -algebra. We say that a positive operator $T \in \mathcal{A}_{FP}(X)$ is *essentially zero* if $\lim_n \frac{Tr(TP_n)}{|VK_n|} = 0$, where P_n is the orthogonal projection onto the space generated by the vertices of K_n in $\ell^2(VX)$.

Proposition 2.2. *Essentially zero operators form the positive part of a closed two-sided ideal $\mathcal{I}(X)$ of $\mathcal{A}_{FP}(X)$.*

Proof. They clearly form a hereditary closed cone. We have to show that such cone is unitary invariant. Indeed, if B has finite propagation r , $P_n B = P_n B P_n(r)$, where $P_n(r)$ denotes the projection on the space generated by the vertices in $\cup_{v \in VK_n} B(v, r)$. Therefore

$$\begin{aligned} Tr(B^* T B P_n) &= Tr(B^* P_n(r) T P_n(r) B P_n) \leq \|B\|^2 Tr(P_n(r) T P_n(r)) \\ &\leq \|B\|^2 Tr(P_n T P_n) + 2\|B\|^2 \|T\| Tr(P_n(r) - P_n). \end{aligned}$$

Moreover, $Tr(P_n(r) - P_n)$ can be estimated by the cardinality of $\cup_{v \in \mathcal{F}K_n} B(v, r)$, hence by $|\mathcal{F}K_n| d^{r+1}$. Regularity of the exhaustion implies $\lim_n \frac{Tr(B^* T B P_n)}{|VK_n|} = 0$, namely $B^* T B$ is essentially zero. Since the cone is closed, we get the invariance for unitaries in $\mathcal{A}_{FP}(X)$. \square

Our first class of amenable graphs is given by the periodic ones. Let Γ be a countable discrete subgroup of automorphisms of X acting freely on X (i.e. any $\gamma \in \Gamma$, $\gamma \neq id$ doesn't have fixed points), and with finite quotient $B := X/\Gamma$. Denote by $F \subset VX$ a set of representatives for VX/Γ , the vertices of the quotient graph B . F is called a *fundamental domain* for the periodic network X .

Let us define a unitary representation of Γ on $\ell^2(VX)$ by $(\lambda(\gamma)f)(x) := f(\gamma^{-1}x)$, for $\gamma \in \Gamma$, $f \in \ell^2(VX)$, $x \in V(X)$. Then the von Neumann algebra $\mathcal{N}(X, \Gamma) := \{\lambda(\gamma) : \gamma \in \Gamma\}'$, of bounded operators on $\ell^2(VX)$ commuting with the action of Γ , inherits a trace given by $Tr_\Gamma(T) = \sum_{x \in F} T(x, x)$, for $T \in \mathcal{N}(X, \Gamma)$. Clearly A , D and Δ belong to $\mathcal{N}(X, \Gamma)$. The following theorem is known, see e.g. [7], Theorem 6.2 for a proof.

¹The definition used here implies $\Delta > 0$, and differs from the standard one adopted in the physics literature.

Theorem 2.3. *Let X be a connected, countably infinite graph, Γ be a countable discrete amenable subgroup of automorphisms of X which acts on X freely and cofinitely. Then X is an amenable graph, and K_n can be chosen in such a way that, for a suitable choice of a sequence $E_n \subset \Gamma$, $V(K_n \setminus \mathcal{F}K_n) \subseteq E_n F \subseteq V(K_n)$, namely K_n is the finite union of copies of F up to $\mathcal{F}K_n$.*

Let us now introduce finite propagation operators, and extend the previously defined trace.

We first notice that, denoting by P_n the projection on $\ell^2(VK_n)$, $Tr_\Gamma(T) = \lim_n \frac{Tr(TP_n)}{|VK_n|}$, for any $T \in \mathcal{N}(X, \Gamma)$. Indeed, by the properties above, the difference between the generic term of the sequence and $Tr_\Gamma(T)$ is infinitesimal.

Next Corollary immediately follows, cf. [13], Corollary 1.5.8.

Corollary 2.4. *Let us denote by $\mathcal{A}(X)$ the space $(\mathcal{N}(X, \Gamma) \cap \mathcal{A}_{FP}(X)) + \mathcal{J}(X)$. Then, $\mathcal{A}(X)$ is a C^* -algebra to which the trace Tr_Γ naturally extends.*

Remark 2.5. All finite rank operators, and hence all compact operators, are essentially zero. The trace Tr_Γ is faithful on $\mathcal{N}(X, \Gamma)$, but it is not faithful on $\mathcal{A}(X)$, in particular it vanishes on compact operators.

Now we discuss small perturbations of amenable periodic graphs. If Y is a finite perturbation of an amenable periodic graph X , namely they only differ for a finite number of points and edges, we can consider both of them as subgraphs of a third graph Z , X and Y being obtained by removing finitely many vertices and finitely many edges from Z . It is not difficult to see that the exhaustion K_n of X can be finitely perturbed to a regular exhaustion K'_n of Z , and that $\mathcal{N}(X, \Gamma)$ is a (possibly non unital) subalgebra of $B(\ell^2(VZ))$.

Reasoning as before, we can consider the unital C^* -algebra $\mathcal{A}(Z) = (\mathcal{N}(X, \Gamma) \cap \mathcal{A}_{FP}(Z)) + \mathcal{J}(Z)$, to which the trace Tr_Γ naturally extends. Since the adjacency operators A_X and A_Y only differ for a finite rank operator, $A_Y \in \mathcal{A}(Z)$, and $Tr_\Gamma(A_X^k) = Tr_\Gamma(A_Y^k)$. More generally, for any continuous function φ on \mathbb{R} , $Tr_\Gamma(\varphi(A_X)) = Tr_\Gamma(\varphi(A_Y))$. This kind of invariance extends to a more general family of small perturbations, which we call density zero perturbations.

For the sake of simplicity, the result below concerns (possibly infinite) perturbations involving only edges. Of course, further finite perturbations can be treated as explained above. The general case of density zero perturbations, studied in [6], can be recovered following the same lines.

Definition 2.6. Let X be an amenable periodic graph, with exhaustion K_n , and consider a graph Y such that $VX = VY$, so that A_X and A_Y both act on the same Hilbert space $\ell^2(VX)$. We say that Y is a *density zero* perturbation of X if $A_X - A_Y$ is essentially zero. In this case, Y is also said to be an *essentially periodic* graph.

Proposition 2.7. *Let X be an amenable periodic graph, with exhaustion K_n , and let Y be a graph with the same vertices as X . Then Y is a density zero perturbation of X if and only if*

$$\lim_n \frac{|(EX \triangle EY) \cap EK_n|}{|VK_n|} = 0,$$

where $EX \triangle EY$ denotes the symmetric difference. In this case, for any continuous function φ on \mathbb{R} ,

$$(2.1) \quad Tr_\Gamma(\varphi(A_X)) = Tr_\Gamma(\varphi(A_Y)).$$

Proof. Clearly $A_X - A_Y$ is essentially zero iff $\lim_n \frac{Tr((A_X - A_Y)^2 P_n)}{|VK_n|} = 0$. A simple calculation shows that $\langle v, (A_X - A_Y)^2 v \rangle = |\{e \in EX \triangle EY : v \in e\}|$. As a consequence, since edges for which both vertices are in K_n should be counted twice,

$$|(EX \triangle EY) \cap EK_n| \leq Tr((A_X - A_Y)^2 P_n) \leq 2|(EX \triangle EY) \cap EK_n|.$$

The thesis follows. Concerning the last equality, setting $T = A_Y - A_X \in \mathcal{I}(X)$, we have $A_Y^n - A_X^n = (A_X + T)^n - A_X^n \in \mathcal{I}(X)$, namely (2.1) holds for $\varphi(t) = t^n$. So the claim is true for any polynomial, and then, using Weierstrass density theorem, for any continuous function. \square

Remark 2.8. We note that for an essentially periodic graph X the C^* -algebra $\mathcal{C}(X)$, and the trace on it, depend in principle on the exhaustion. However, previous Proposition implies that on geometric operators, such as the adjacency A and its continuous functional calculi, the value of the trace is uniquely determined.

3. STATISTICAL MECHANICS ON AMENABLE GRAPHS

The main aim of the present paper is to investigate in full generality the thermodynamics of free Bosons (Baarden–Cooper pairs) on inhomogeneous networks with pure hopping Hamiltonian (i.e. the opposite of the adjacency matrix on the graph). Thus for the convenience of the reader, we report some standard notions useful in the sequel.

Let (\mathfrak{A}, α) be a dynamical system consisting of a (noncommutative) C^* -algebra and a one parameter group of $*$ -automorphism α . The state $\omega \in \mathcal{S}(\mathfrak{A})$ satisfies the *KMS boundary condition* at inverse temperature β , which we suppose to be always different from zero, if

- (i) $t \mapsto \omega(A\alpha_t(B))$ is a continuous function for every $A, B \in \mathfrak{A}$,
- (ii) $\int \omega(A\alpha_t(B))f(t)dt = \int \omega(\alpha_t(B)A)f(t + i\beta)dt$ whenever $f \in \widehat{\mathcal{D}}$, where “ $\widehat{}$ ” stands for the Fourier transform.

Here, \mathcal{D} is the space of smooth compactly supported functions on \mathbb{R} .

The C^* -algebras considered here are those arising from the Canonical Commutation Relations (CCR for short). Namely, let \mathfrak{h} be a pre-Hilbert space and consider the following (formal) relations between the annihilators $a(f)$, and creators $a^+(g)$, $f, g \in \mathfrak{h}$

$$(3.1) \quad a(f)a^+(g) - a^+(g)a(f) = \langle f, g \rangle.$$

It is well-known that the relations (3.1) cannot be realized by bounded operators. A standard way to realize (3.1) is to look at the symmetric Fock space $\mathcal{F}_+(\mathfrak{h})$ on which the annihilators and creators naturally act as unbounded closed (mutually adjoint) operators. This concrete representation of the CCR is called *the Fock representation*.

An equivalent description for the CCR is to put $\Phi(f) := \overline{a(f) + a^+(f)}/\sqrt{2}$, and define the *Weyl operators* $W(f) := \exp i\Phi(f)$. The Weyl operators are unitary and satisfy the rule

$$(3.2) \quad W(f)W(g) = e^{-i\frac{Im(f,g)}{2}} W(f+g), \quad f, g \in \mathfrak{h}.$$

The CCR algebra $CCR(\mathfrak{h})$ is precisely the C^* -algebra generated by $\{W(f)\}_{f \in \mathfrak{h}}$.

Let H be a positive operator acting on \mathfrak{h} , and suppose that $e^{itH}\mathfrak{h} \subset \mathfrak{h}$. Then the one-parameter group of Bogoliubov automorphisms $T_t f := e^{itH} f$ defines a

one-parameter group of $*$ -automorphisms α_t of $\text{CCR}(\mathfrak{h})$ by putting $\alpha_t(W(f)) := W(e^{itH}f)$.

A representation π of the CCR algebra $\text{CCR}(\mathfrak{h})$ is *regular* if the unitary group $t \in \mathbb{R} \mapsto \pi(W(tf))$ is continuous in the strong operator topology, for any $f \in \mathfrak{h}$. A state φ on $\text{CCR}(\mathfrak{h})$ is regular if the associated GNS representation is regular. It simply means that the functions $\{\varphi(W(tf))\}$ are continuous, for any $f \in \mathfrak{h}$. The *quasi-free* states of CCR algebras are of interest for our purposes. They are analytic states ω uniquely determined by the two-point functions $\omega(a^+(f)a(g))$, $f, g \in \mathfrak{h}$.

Let X be an infinite graph, \mathfrak{h} a subspace of $\ell^2(X)$, which contains the indicator functions of all finite subregions Λ of the graph X . A representation π of the CCR algebra $\text{CCR}(\mathfrak{h})$ is said to be *locally normal* (w.r.t. the Fock representation) if $\pi|_{\text{CCR}(\ell^2(\Lambda))}$ is quasi-equivalent to the Fock representation of $\text{CCR}(\ell^2(\Lambda))$. A state on $\text{CCR}(\mathfrak{h})$ is locally normal if the associated GNS representation is. A locally normal state φ does have finite local density

$$\rho_\Lambda(\varphi) := \frac{1}{|\Lambda|} \sum_{j \in \Lambda} \varphi(a^+(\delta_j)a(\delta_j))$$

even if the mean density might be infinite (e.g. if $\lim_{\Lambda \uparrow X} \rho_\Lambda(\varphi) = +\infty$).

Lemma 3.1. *If φ is locally normal, then $\varphi|_{\text{CCR}(\ell^2(\Lambda))}$ is regular for any finite subregion Λ .*

Proof. Since $\pi_\varphi|_{\text{CCR}(\ell^2(\Lambda))}$ is quasi equivalent to the Fock representation of $\text{CCR}(\ell^2(\Lambda))$, they are unitary equivalent up to multiplicity (cf. thm 2.4.26 [2]). The result follows since the Fock representation is regular. \square

We now specialize to the following situation. Let $\Lambda_n \uparrow X$ be a sequence of finite regions invading the graph X , together with a sequence of states $\{\omega_{\Lambda_n}\}$ on $\text{CCR}(\ell^2(\Lambda_n))$ such that the following limit

$$\lim_n \omega_{\Lambda_n}(a^+(v)a(v)) =: q(v)$$

exists (possibly $+\infty$) for each $v \in \mathfrak{h}$.

Lemma 3.2. *Suppose that*

$$(3.3) \quad \lim_n \omega_{\Lambda_n}(a^+(\delta_j)a(\delta_j)) = +\infty$$

for some $j \in X$. Then $\omega(W(v)) := \lim_n \omega_{\Lambda_n}(W(v))$ does not define any locally normal state on $\text{CCR}(\mathfrak{h})$, where $\mathfrak{h} \subset \ell^2(X)$ is any subspace containing the finite supported sequences.

Proof. Let j be contained in the finite region Λ . By Lemma 3.1, it is enough to show that $\omega|_{\mathfrak{A}_\Lambda}$ is not regular. We have (cf. [3], Example 5.2.18)

$$\omega(W(\lambda\delta_j)) = \lim_n \omega_{\Lambda_n}(W(\lambda\delta_j)) = e^{-\frac{1}{4}\lambda^2} \lim_n \exp\left(-\frac{\lambda^2}{2} \omega_{\Lambda_n}(a^+(\delta_j)a(\delta_j))\right) = 0.$$

The thesis follows as ω cannot be regular. \square

For equivalent characterizations of the KMS boundary condition and general results on the CCR the reader is referred to [3] and the references cited therein.

In this work, we consider amenable graphs (also called amenable networks) which are essentially-periodic, namely finite or density-zero perturbations of periodic graphs as described above. We assume a regular exhaustion $\{\Lambda_n\}$ is given,

and denote by τ the canonical trace on $\mathcal{A}(X)$. We also denote with an abuse of notation, $\ell^2(X) \equiv \ell^2(VX)$ for the network X . In this section we shall introduce the main thermodynamic properties and quantities.

Fix a positive operator $H \in \mathcal{A}(X)$ (the Hamiltonian) and denote by N_H its *integrated density of states*, (see e.g. [12]) that is $N_H(\lambda) := \tau(E[0, \lambda])$, where $H = \int \lambda dE(\lambda)$ is the spectral decomposition of H .

Let $\{H_{\Lambda_n}\}$ be its finite volume truncation w.r.t. the exhaustion $\{\Lambda_n\}$, i.e. $H_{\Lambda_n} := P_n H P_n$, and denote by $N_{H_{\Lambda_n}}(\lambda) := \frac{1}{|\Lambda_n|} \text{Tr}(E_{H_{\Lambda_n}}[0, \lambda]) \equiv \frac{1}{|\Lambda_n|} |\{t \in \sigma(H_{\Lambda_n}) : t \leq \lambda\}|$. Define

$$(3.4) \quad \begin{aligned} E_0(H) &:= \lim_{\Lambda_n \uparrow X} \left(\inf \text{supp} (N_{H_{\Lambda_n}}) \right) \\ E_m(H) &:= \inf \text{supp} \left(\lim_{\Lambda_n \uparrow X} N_{H_{\Lambda_n}} \right) \equiv \inf \text{supp} (N_H). \end{aligned}$$

Remark 3.3.

- (i) The limit in (3.4) exists as a consequence of Lemma 3.4.
- (ii) Because of Proposition 3.5, we have $\lim_{\Lambda_n \uparrow X} N_{H_{\Lambda_n}} = N_H$, as distribution functions.
- (iii) If A is the adjacency matrix of X , and $H := \|A\| - A$, then $E_0(H) = \lim_{\Lambda_n \uparrow X} \|A\| - \|A_{\Lambda_n}\| = 0$, while $E_m(H) = \|A\| - \|\pi_\tau(A)\|$, where π_τ is the GNS representation induced by τ . Moreover, if X is a periodic graph, then it is shown in Theorem 5.2 that $E_m(H) = 0$.

Obviously, $E_0(H) \leq E_m(H)$. If $E_0(H) < E_m(H)$ we say that there is a (*low energy*) *hidden spectrum*, see e.g. [4]. For the infinite graphs below, we shall always assume $E_0(H) = 0$.

Lemma 3.4. *Let X be a countable graph, $H \in \mathcal{B}(\ell^2(X))$ a positive operator, and let $\Lambda_1 \subset \Lambda_2$ be finite subgraphs of X . Then*

$$\min \sigma(H_{\Lambda_1}) \geq \min \sigma(H_{\Lambda_2}).$$

Proof. Let λ be the minimum eigenvalue of H_{Λ_1} , and v the relative normalised eigenvector. Since $v \in \ell^2(\Lambda_1)$, we have $\langle v, H_{\Lambda_2} v \rangle = \langle v, H_{\Lambda_1} v \rangle = \lambda$, that is $\lambda \in \{\langle w, H_{\Lambda_2} w \rangle : w \in \ell^2(\Lambda_2), \|w\| = 1\}$. Observe that this set is contained in $\text{conv}(\sigma(H_{\Lambda_2}))$, since, if $H_{\Lambda_2} = \int \lambda dE(\lambda)$ is the spectral decomposition, we have $\langle v, H_{\Lambda_2} v \rangle = \int \lambda d\langle v, E(\lambda) v \rangle$, and the claim follows from the fact that $d\langle v, E(\lambda) v \rangle$ is a probability measure on $\sigma(H_{\Lambda_2})$. Finally, $\min \sigma(H_{\Lambda_2}) = \min \text{conv} \sigma(H_{\Lambda_2}) \leq \lambda = \min \sigma(H_{\Lambda_1})$. \square

Proposition 3.5. *Let X be an essentially periodic graph with regular exhaustion $\{\Lambda_n\}$, $H \in \mathcal{A}(X)$ a positive operator. Then, for any continuous function $\varphi : [0, \infty) \rightarrow \mathbb{C}$, we have*

$$(i) \quad \lim_{n \rightarrow \infty} \frac{\text{Tr}(\varphi(H_{\Lambda_n}))}{|\Lambda_n|} = \tau(\varphi(H)),$$

$$(ii) \quad \int \varphi dN_H = \lim_{\Lambda_n \uparrow X} \int \varphi dN_{H_{\Lambda_n}}.$$

Proof. Let us denote by E_n the orthogonal projection from $\ell^2(X)$ onto $\ell^2(\Lambda_n)$.

Observe that, for $k \geq 2$,

$$\begin{aligned} \text{Tr}(E_n H^k E_n) &= \text{Tr}(E_n (H(E_n + E_n^\perp))^k E_n) \\ &= \text{Tr}((E_n H E_n)^k) + \sum_{\substack{\sigma \in \{-1,1\}^{k-1} \\ \sigma \neq \{1,1,\dots,1\}}} \text{Tr}(E_n \prod_{j=1}^{k-1} (H E_n^{\sigma_j}) H E_n), \end{aligned}$$

where E_n^{-1} stands for E_n^\perp , and

$$|\text{Tr}(E_n \prod_{j=1}^{k-1} (H E_n^{\sigma_j}) H E_n)| = |\text{Tr}(\dots E_n H E_n^\perp \dots)| \leq \|H\|^{k-1} \text{Tr}(|E_n H E_n^\perp|).$$

Now assume H has propagation r . Then,

$$E_n H E_n^\perp = E_n E(B_r(V\Lambda_n^c)) H E_n^\perp = E(\Lambda_n \cap B_r(V\Lambda_n^c)) H E_n^\perp,$$

hence

$$\text{Tr}(|E_n H E_n^\perp|) \leq \|H\| |B_{r-1}(\mathcal{F}\Lambda_n)| \leq \|H\| |\mathcal{F}\Lambda_n| d^{k-1},$$

where d is the maximal degree of X . As a consequence we obtain

$$\frac{1}{|\Lambda_n|} |\text{Tr}(E_n H^k E_n) - \text{Tr}((E_n H E_n)^k)| \leq (2d)^{k-1} \|H\|^k \frac{|\mathcal{F}\Lambda_n|}{|\Lambda_n|} \rightarrow 0.$$

Setting $\tau_n = \frac{1}{|\Lambda_n|} \text{Tr}$, we have proved that $\lim_n \tau_n(E_n p(H) E_n - p(E_n H E_n)) = 0$ for any positive, finite propagation operator H in $\mathcal{A}(X)$. Since $|\tau_n(A)| \leq \|A\|$, the result follows by Weierstrass density theorem and the definition of $\mathcal{A}(X)$. \square

Definition 3.6. Let H be a positive operator in $\mathcal{A}(X)$. Then, for any *inverse temperature* $\beta > 0$, and for any *chemical potential* $\mu \leq 0$, we define the *density* of H as

$$(3.5) \quad \rho_H(\beta, \mu) := \int \frac{dN_H(h)}{e^{\beta(h-\mu)} - 1},$$

and the *critical density* of H as

$$(3.6) \quad \rho_c^H(\beta) := \int \frac{dN_H(h)}{e^{\beta h} - 1} \equiv \rho_H(\beta, 0).$$

Let us recall that H is called *recurrent* if the matrix elements $(\delta_x, H^{-1}\delta_x)$ are infinite, and *transient* if the matrix elements $(\delta_x, H^{-1}\delta_x)$ are finite.

We say that BEC takes place for a given equilibrium state if a suitable portion of the particles occupies the lowest energy state.

Remark 3.7.

- (i) It is also customary to fix the *activity* $z := e^{\beta\mu}$, instead of the chemical potential.
- (ii) If we choose H as the graph Laplacian, transience and recurrence have a probabilistic interpretation. In fact, $(\delta_x, \Delta^{-1}\delta_y) = \sum_{n=0}^{\infty} p_n(x, y)$, where $p_n(x, y)$ is the probability of passing from x to y in n steps, and the transition probability $p_1(x, y)$ is set to $(\deg x)^{-1}$ if x and y are adjacent, and to 0 otherwise. As a consequence, if the graph is connected, $(\delta_x, \Delta^{-1}\delta_x)$ is finite for all x if and only if it is finite for a single x . Then transience corresponds to the following property of a random walk on X : the random walk starting at a point x returns almost surely to x infinitely many times. Conversely, a random walk is recurrent if the probability of starting at x and returning to x infinitely many times is zero. Interpreting a graph as an

electrical network, transience means that the resistance between a point and infinity is finite. For further results on transience and recurrence see e.g. [16]

(iii) For the standard homogeneous models investigated in literature (i.e. the statistical mechanics of free Bosons on \mathbb{R}^d (cf. [3]), or on lattices with period \mathbb{Z}^d (cf. [10] and the references cited therein)),² it is well-known that there exist equilibrium states exhibiting BEC if $\rho_c^H(\beta) < +\infty$. This is also known to be equivalent to the transience of the graph, or to the fact that the growth of the graph is greater than 2. As we shall see in the following sections, for the nonhomogeneous models treated in the present paper new phenomena (as for example the lack of the local normality of the resulting state in the thermodynamical limit) can happen.

(iv) Since we assumed H to be bounded, the critical density is finite, namely the integral (3.6) converges, iff

$$(3.7) \quad \int_{[0, \|H\|]} \frac{dN_H(h)}{h} < +\infty.$$

In particular, hidden spectrum implies finite critical density. A large class of examples of essentially periodic graphs exhibiting hidden spectrum will be described below.

Now we recall how equilibrium states are usually constructed. Given a positive Hamiltonian H as above on the essentially periodic graph X , one fixes an inverse temperature $\beta > 0$ and a density of particles ρ , and determines the chemical potential $\mu(\Lambda_n)$ such that

$$\rho_{H_{\Lambda_n}}(\beta, \mu(\Lambda_n)) = \rho.$$

To simplify the exposition, we suppose also that $\mu(\Lambda_n) \rightarrow \mu$, for some $\mu \in \mathbb{R}$, possibly by passing to a subsequence. We necessarily have $\mu(\Lambda_n) < E_0(H_{\Lambda_n})$. Since $E_0(H_{\Lambda_n}) \rightarrow E_0(H) = 0$, we get $\mu \leq 0$.

The finite volume state with density ρ describing the Gibbs grand canonical ensemble in the volume Λ_n can be defined by the two-point function

$$(3.8) \quad \omega_{\Lambda_n}(a^+(\xi)a(\eta)) = \langle \eta, (e^{\beta(H_{\Lambda_n} - \mu(\Lambda_n)I)} - 1)^{-1} \xi \rangle,$$

where $\xi, \eta \in \ell^2(X)$. Thermodynamical states are then described as limits of the finite volume states above.

Let us now study the behavior of the density in the infinite volume limit.

Proposition 3.8. *Let X be an essentially periodic graph with regular exhaustion $\{\Lambda_n\}$, with positive Hamiltonian $H \in \mathcal{A}(X)$, and assume $N_H(0) = 0$. Then, we have*

$$\lim_{\varepsilon \downarrow 0} \lim_{\Lambda_n \uparrow X} \int \frac{f_\varepsilon(h)}{e^{\beta(h - \mu(\Lambda_n))} - 1} dN_{H_{\Lambda_n}}(h) = \rho_H(\beta, \mu),$$

where f_ε is the continuous mollifier

$$f_\varepsilon(x) = \begin{cases} 0, & 0 \leq x \leq \varepsilon \\ \frac{x - \varepsilon}{\varepsilon}, & \varepsilon < x \leq 2\varepsilon \\ 1, & 2\varepsilon < x. \end{cases}$$

²As noticed in [4], the BEC behaviour of \mathbb{Z}^d -lattices only depends on d . This can be seen as the dispersion law $\mathbf{p} \mapsto \varepsilon_H(\mathbf{p})$ of a periodic Schrödinger operator $H := \Delta + V$ on \mathbb{R}^d or a lattice, has the same asymptotics near $\mathbf{0}$ as that of the Laplacian Δ , see e.g. [15].

Proof. We compute

$$\begin{aligned} & \left| \int \frac{f_\varepsilon(h)}{e^{\beta(h-\mu(\Lambda_n))} - 1} dN_{H_{\Lambda_n}}(h) - \int \frac{f_\varepsilon(h)}{e^{\beta(h-\mu)} - 1} dN_H(h) \right| \\ & \leq \left| \int \left(\frac{1}{e^{\beta(h-\mu(\Lambda_n))} - 1} - \frac{1}{e^{\beta(h-\mu)} - 1} \right) f_\varepsilon(h) dN_{H_{\Lambda_n}}(h) \right| \\ & + \left| \int \frac{f_\varepsilon(h)}{e^{\beta(h-\mu)} - 1} dN_{H_{\Lambda_n}}(h) - \int \frac{f_\varepsilon(h)}{e^{\beta(h-\mu)} - 1} dN_H(h) \right| \rightarrow 0, \end{aligned}$$

since $\frac{1}{e^{\beta(h-\mu(\Lambda_n))} - 1} \rightarrow \frac{1}{e^{\beta(h-\mu)} - 1}$, uniformly on the support of f_ε , and the second summand goes to zero because the finite volume sequences of traces converge to the infinite volume trace for fixed traceable operators, as follows from Proposition 3.5 (ii). The proof follows as $f_\varepsilon \nearrow 1$ whenever $\varepsilon \searrow 0$, essentially everywhere by taking into account that we assumed $N_H(0) = 0$.³ \square

Let us observe that the quantity

$$n_0 = n_0(\beta, \rho) := \lim_{\varepsilon \downarrow 0} \lim_{\Lambda_n \uparrow X} \int \frac{1 - f_\varepsilon(h)}{e^{\beta(h-\mu(\Lambda_n))} - 1} dN_{H_{\Lambda_n}}(h)$$

is well-defined and independent of the particular choice of the mollifier f_ε whenever the last converges monotonically to $\chi_{(0,+\infty)}$. We have

$$(3.9) \quad \rho = n_0 + \rho(\beta, \mu).$$

From now on, we suppose that the Hamiltonians H , H_Λ are those based on the adjacency matrix if it is not otherwise specified, and drop some subscripts.

Proposition 3.9. *Let X be an essentially periodic graph, with adjacency matrix A and $H = \|A\| - A$. Then $N_H(0) = 0$, i.e. the integrated density of states is continuous in zero.*

Proof. Let us assume first that X is a periodic amenable graph. Then, if $N_H(0) \neq 0$, the 0-eigenspace of H is non-trivial, namely there is an ℓ^2 Perron-Frobenius eigenvector for A , which is necessarily unique. However, since H is G invariant, such vector is also periodic. This is absurd.

Let now X be essentially periodic. If there is no hidden spectrum, then $N_H(\lambda)$ is the same as that of a periodic graph, hence the result follows from the previous case. On the other hand, hidden spectrum immediately implies $N_H(0) = 0$. \square

Lemma 3.10. *Let X be an essentially periodic graph with regular exhaustion $\{\Lambda_n\}$, A the adjacency matrix of X . Let $\{\mu_n\} \subset (-\infty, 0]$, $\mu_n \rightarrow \mu$, and set $H := (\|A\| - \mu)I - A$, and $H_n := (\|A\| - \mu_n)I - A_{\Lambda_n}$. Let $\varphi : [0, \infty) \rightarrow \mathbb{C}$ be a continuous function. Then, for any $\xi \in \ell^2(X)$, $\varphi(H_n)\xi \rightarrow \varphi(H)\xi$.*

Proof. Let $x \in VX$; then, for $n \in \mathbb{N}$ large enough, $A_{\Lambda_n}\delta_x = A\delta_x$, so that $A_{\Lambda_n}\xi = A\xi$, for $\xi \in \ell^2(X)$ and with finite support. Let $\xi \in \ell^2(X)$ be arbitrary, and, for any $\varepsilon > 0$, let $\xi_\varepsilon \in \ell^2(X)$ and with finite support be such that $\|\xi - \xi_\varepsilon\| < \varepsilon$. Then, for some $n_\varepsilon \in \mathbb{N}$, we have $A_{\Lambda_n}\xi_\varepsilon = A\xi_\varepsilon$, for any $n > n_\varepsilon$. Therefore, for $n > n_\varepsilon$ we have $\|A_{\Lambda_n}\xi - A\xi\| \leq \|A_{\Lambda_n}(\xi - \xi_\varepsilon)\| + \|A_{\Lambda_n}\xi_\varepsilon - A\xi_\varepsilon\| + \|A(\xi - \xi_\varepsilon)\| \leq 2\|A\|\|\xi - \xi_\varepsilon\| < 2\|A\|\varepsilon$. Hence $H_n\xi \rightarrow H\xi$.

³Recall that $\mu = 0$ is allowed and describes the most interesting situation of the BEC regime, see below.

Therefore, for any $k \in \mathbb{N}$, $H_n^k \xi \rightarrow H^k \xi$, and the claim is true for any polynomial. Finally, let $\varphi : [0, \infty) \rightarrow \mathbb{C}$ be a continuous function, and, for any $\varepsilon > 0$, let p be a polynomial such that $\|\varphi - p\|_\infty < \varepsilon$. Then, for any $\xi \in \ell^2(X)$, we have $\|\varphi(H_n)\xi - \varphi(H)\xi\| \leq \|\varphi(H_n)\xi - p(H_n)\xi\| + \|p(H_n)\xi - p(H)\xi\| + \|p(H)\xi - \varphi(H)\xi\| \leq 2\|\varphi - p\|_\infty \|\xi\| + \|p(H_n)\xi - p(H)\xi\|$, from which the claim follows. \square

Theorem 3.11. *Let X be an essentially periodic graph with regular exhaustion $\{\Lambda_n\}$, A the adjacency matrix of X , $H := \|A\| - A$. Then*

- (i) $\mu < 0$ if and only if $\rho < \rho_c(\beta)$,
- (ii) for any $\beta > 0$, $\mu < 0$, the sequence (3.8) converges pointwise to a state ω , whose two-point function is given by

$$\omega(a^+(\xi)a(\eta)) = \langle \eta, (e^{\beta(H-\mu I)} - 1)^{-1} \xi \rangle.$$

Moreover, the density $\rho(\omega)$ of the state ω , defined by

$$(3.10) \quad \rho(\omega) := \lim_{\Lambda_n \uparrow X} \frac{1}{|\Lambda_n|} \sum_{j \in \Lambda_n} \omega(a^+(\delta_j)a(\delta_j)),$$

satisfies $\rho(\omega) = \rho(\beta, \mu)$.

- (iii) The transience of A is a necessary condition for the existence of locally normal states on $\text{CCR}(\mathfrak{h})$ at or above the critical density, i.e. for $\mu = 0$. Here $\mathfrak{h} \subset \ell^2(X)$ is a subspace containing the functions with finite support.

Proof. (i) (\implies) If $\mu(\Lambda_n) \rightarrow \mu < 0$, we can suppose that $z(\Lambda_n) \equiv e^{\beta\mu(\Lambda_n)} \leq K < 1$ for each n . We have

$$\begin{aligned} \int \frac{1 - f_\varepsilon(h)}{e^{\beta(h-\mu(\Lambda_n))} - 1} dN_{H_{\Lambda_n}}(h) &\leq \frac{z(\Lambda_n)}{1 - z(\Lambda_n)} \int (1 - f_\varepsilon(h)) dN_{H_{\Lambda_n}}(h) \\ &\leq \frac{K}{1 - K} \int (1 - f_\varepsilon(h)) dN_{H_{\Lambda_n}}(h) \rightarrow \frac{K}{1 - K} N_H(0) = 0, \end{aligned}$$

whenever $\Lambda_n \uparrow X$, and then $\varepsilon \searrow 0$. Therefore, $n_0 = 0$, so that $\rho = \rho(\beta, \mu) < \rho(\beta, 0) \equiv \rho_c(\beta)$.

(\impliedby) Conversely, suppose that $\rho < \rho_c(\beta)$ and $\mu(\Lambda_n) \rightarrow 0$. Then

$$\begin{aligned} \rho_c(\beta) > \rho &= \int \frac{1 - f_\varepsilon(h)}{e^{\beta(h-\mu(\Lambda_n))} - 1} dN_{H_{\Lambda_n}}(h) + \int \frac{f_\varepsilon(h)}{e^{\beta(h-\mu(\Lambda_n))} - 1} dN_{H_{\Lambda_n}}(h) \\ &\geq \int \frac{f_\varepsilon(h)}{e^{\beta(h-\mu(\Lambda_n))} - 1} dN_{H_{\Lambda_n}}(h) \rightarrow \rho(\beta, 0) \equiv \rho_c(\beta), \end{aligned}$$

whenever $\Lambda_n \uparrow X$, and then $\varepsilon \searrow 0$. As this is impossible, we must have $\mu < 0$.

(ii) The fact that $\omega_{\Lambda_n} \rightarrow \omega$ in the $*$ -weak topology follows from $H_{\Lambda_n} \rightarrow H$ in the strong operator topology. Indeed, since $H_n - \mu(\Lambda_n) \rightarrow H - \mu$ strongly, it follows from Lemma 3.10 that $(e^{\beta(H_n - \mu(\Lambda_n))} - 1)^{-1} \rightarrow (e^{\beta(H - \mu)} - 1)^{-1}$ strongly [because $\mu < 0$ so that $1 \notin \sigma(e^{\beta(H - \mu)})$]. Therefore, for any $\xi, \eta \in \ell^2(X)$, we have $\omega_{\Lambda_n}(a^+(\xi)a(\eta)) = \langle \eta, (e^{\beta(H_n - \mu(\Lambda_n))} - 1)^{-1} \xi \rangle \rightarrow \langle \eta, (e^{\beta(H - \mu)} - 1)^{-1} \xi \rangle = \omega(a^+(\xi)a(\eta))$. Since the vector space $V := \text{span} \{a^+(\xi)a(\eta) : \xi, \eta \in \ell^2(X)\}$ is dense in CCR , then ω_{Λ_n} is a Cauchy sequence in the weak* topology, so it converges. Indeed, for any $T \in \text{CCR}$, $\varepsilon > 0$, there is $T_\varepsilon \in V$ such that $\|T - T_\varepsilon\| < \varepsilon$, so that, for any m, n large enough, $|\omega_{\Lambda_m}(T) - \omega_{\Lambda_n}(T)| \leq |\omega_{\Lambda_m}(T) - \omega_{\Lambda_m}(T_\varepsilon)| + |\omega_{\Lambda_m}(T_\varepsilon) - \omega_{\Lambda_n}(T_\varepsilon)| + |\omega_{\Lambda_n}(T_\varepsilon) - \omega_{\Lambda_n}(T)| \leq 3\varepsilon$.

Finally,

$$\int \frac{dN_{H_{\Lambda_n}}(h)}{e^{\beta(h-\mu(\Lambda_n))} - 1} \rightarrow \int \frac{dN_H(h)}{e^{\beta(h-\mu)} - 1}$$

as in the proof of Proposition 3.8, because

$$\frac{1}{e^{\beta(h-\mu(\Lambda_n))} - 1} \rightarrow \frac{1}{e^{\beta(h-\mu)} - 1},$$

uniformly on $[0, \|H\|]$, and the measures induced by the $N_{H_{\Lambda_n}}$ converge in the $*$ -weak topology to that induced by N_H .

(iii) Indeed, the following limit should be finite:

$$\begin{aligned} & \lim_{\varepsilon \downarrow 0} \lim_{\Lambda_n \uparrow X} \langle \xi, f_\varepsilon(H_n)(e^{\beta(H_n - \mu_n)} - 1)^{-1} \eta \rangle \\ &= \lim_{\varepsilon \downarrow 0} \lim_{\Lambda_n \uparrow X} \langle \xi, f_\varepsilon(H_n)(e^{\beta H_n} - 1)^{-1} \eta \rangle \\ &+ \lim_{\varepsilon \downarrow 0} \lim_{\Lambda_n \uparrow X} \langle \xi, f_\varepsilon(H_n)((e^{\beta(H_n - \mu_n)} - 1)^{-1} - (e^{\beta H_n} - 1)^{-1}) \eta \rangle. \end{aligned}$$

Since for $\mu_n \rightarrow 0$ the second summand is zero, we need the finiteness of the first. By Proposition 3.10,

$$\begin{aligned} & \lim_{\varepsilon \downarrow 0} \lim_{\Lambda_n \uparrow X} \langle \xi, f_\varepsilon(H_n)(e^{\beta H_n} - 1)^{-1} \eta \rangle \\ &= \lim_{\varepsilon \downarrow 0} \langle \xi, f_\varepsilon(H)(e^{\beta H} - 1)^{-1} \eta \rangle \\ &= \langle \xi, (e^{\beta H} - 1)^{-1} \eta \rangle. \end{aligned}$$

The finiteness of $\langle \xi, (e^{\beta H} - 1)^{-1} \eta \rangle$ when ξ, η have finite support is exactly the transience of A_X . \square

Remark 3.12.

(i) Let us observe that $n_0 > 0$ if and only if $\rho > \rho_c(\beta)$. Indeed, by definition, $n_0 = 0$ whenever $\mu < 0$, hence $n_0 > 0 \Rightarrow \mu = 0$. As a consequence, $0 < n_0 = \rho - \rho(\beta, \mu) = \rho - \rho_c(\beta)$, i.e. $\rho > \rho_c(\beta)$. Conversely, $\rho > \rho_c(\beta)$ implies $\mu = 0$, hence $n_0 = \rho - \rho(\beta, \mu) = \rho - \rho_c(\beta) > 0$.

(ii) Observe that $n_0 > 0$ can be obtained only if $\mu(\Lambda_n) \rightarrow 0$. Indeed, $\mu = 0$ is a necessary condition for the occurrence of BEC.

4. SOME RESULTS ON PERRON-FROBENIUS EIGENVECTORS

Let X be a finite connected graph, A its adjacency matrix. By Perron-Frobenius Theorem there exists a unique eigenvector with eigenvalue $\|A\|$, and it is the unique eigenvector having strictly positive entries. When X is infinite, the existence of a square summable Perron-Frobenius eigenvector is no longer guaranteed. If such vector exists is unique and has strictly positive entries. However, if X has bounded degree, the equation $Av = \|A\|v$ makes sense also for vectors which are simply functions $v : X \mapsto \mathbb{C}$. Vectors of this kind satisfying the equation $Av = \|A\|v$ will be called generalized Perron-Frobenius eigenvectors. A generalized Perron-Frobenius eigenvector has strictly positive entries but it is not necessarily unique. Indeed it is unique if the graph is A -recurrent, or if it is tree-like [16, 14]. Now we show the existence of such (generalized) Perron-Frobenius eigenvectors for the network under consideration.

Assume Λ_n is an exhaustion for X , namely an increasing family of connected finite subgraphs whose union is X , and choose a vertex $x_0 \in V\Lambda_1$. Then let v_n be

the Perron-Frobenius vector for $A_n := A_{\Lambda_n}$ normalized by $\langle \delta_{x_0}, v_n \rangle = 1$. We extend all these vectors to zero outside Λ_n .

Proposition 4.1. *With the above notation, A has a generalized Perron–Frobenius eigenvector.*

Proof. We first show that $\|A\| = \lim_n \|A_n\|$. Indeed, if $Y \subset Z$ is a proper inclusion of graphs with Y finite and Z connected, and w is the norm-one Perron-Frobenius vector for A_Y , then $\|A_Y\| = \langle w, A_Y w \rangle < \langle w, A_Z w \rangle \leq \|A_Z\|$. Therefore the sequence $\|A_n\|$ is strictly increasing and bounded by $\|A\|$. Assume now ad absurdum that $\lim_n \|A_n\| < \|A\|$. Then we could find a norm-one vector z with finite support such that $\langle z, Az \rangle > \lim_n \|A_n\|$. Choosing n large enough, we get $\langle z, Az \rangle = \langle z, A_n z \rangle \leq \|A_n\|$, which gives a contradiction.

Finally we construct a generalized Perron-Frobenius vector for X . Since

$$1 = \langle \delta_{x_0}, v_n \rangle = \|A_n\|^{-1} \langle \delta_{x_0}, A_n v_n \rangle = \|A_n\|^{-1} \sum_{x \sim x_0} \langle \delta_x, v_n \rangle \geq \|A_n\|^{-1} \max_{x \sim x_0} \langle \delta_x, v_n \rangle,$$

then $x \sim x_0$ implies $\langle \delta_x, v_n \rangle \leq \|A_n\|$. By induction we get $d(x, x_0) \leq d$ implies $\langle \delta_x, v_n \rangle \leq \|A_n\|^d$. From this we can obtain a subsequence v_{n_k} such that $\langle \delta_x, v_{n_k} \rangle$ converges for any $x \in X$. Let us denote by v the vector s.t. $\lim_k \langle \delta_x, v_{n_k} \rangle = \langle \delta_x, v \rangle$. We observe that, given $x \in X$, for sufficiently large k , we have $A\delta_x = A_{n_k}\delta_x$. Therefore

$$\langle \delta_x, Av \rangle = \langle A\delta_x, v \rangle = \lim_k \langle \delta_x, A_{n_k} v_{n_k} \rangle = \lim_k \|A_{n_k}\| \langle \delta_x, v_{n_k} \rangle = \langle \delta_x, \|A\|v \rangle,$$

which means that v is a generalized Perron-Frobenius vector for X . \square

5. NO HIDDEN SPECTRUM FOR PERIODIC GRAPHS

In this section we want to show that there is no hidden spectrum for the adjacency operator on a periodic graph. This means that the bottom of the spectrum for the energy operator $\|A\| - A$ coincides with the infimum of the support of the spectral measure of the energy operator in the trace representation, or, equivalently, that the supremum of the spectrum of A coincides with the supremum of the spectrum for $\pi(A)$, where π is the trace representation.

We have already proved that, given an exhaustion Λ_n of the graph X , and denoting with A_n the adjacency operator for Λ_n , then $\|A_n\| \nearrow \|A\|$.

Let now $v \in \ell^2(VX)$ be a unit vector with support contained in Λ_n , $\gamma \in \Gamma$ and consider the projection operator $P_v(\gamma)$ on the vector $\lambda(\gamma)v$. Let us observe that, since $P_v(\gamma)$ is a projection operator on $\lambda(\gamma)v$, for any bounded operator C ,

$$\text{Tr}(CP_v(\gamma)) = \langle \lambda(\gamma)v, C\lambda(\gamma)v \rangle.$$

We need the following.

Lemma 5.1. *Let X be a periodic graph, K a finite subgraph, v the normalised Perron-Frobenius eigenvector of K , and consider, for any $\gamma \in \Gamma$, the projection operator $P_v(\gamma)$ on the vector $\lambda(\gamma)v$. Then the series $\sum_{\gamma \in \Gamma} P_v(\gamma)$ converges strongly to an operator T which belongs to the von Neumann algebra $\mathcal{N}(X, \Gamma)$.*

Proof. We need to show that, for any $\varepsilon > 0$, and $w = \sum_{y \in X} w_y \delta_y \in \ell^2(X)$, there is a finite set $F_\varepsilon \subset \Gamma$ such that, for all finite sets $E \subset \Gamma \setminus F_\varepsilon$, we have $\|\sum_{\gamma \in E} P(\gamma)w\|^2 < \varepsilon$.

Indeed, for any finite $E \subset \Gamma$, we have

$$\begin{aligned}
\left\| \sum_{\gamma \in E} P_v(\gamma) w \right\|^2 &= \sum_{x \in X} \left| \langle \delta_x, \sum_{\gamma \in E} P_v(\gamma) w \rangle \right|^2 \\
&= \sum_{x \in X} \left| \sum_{\gamma \in E} \sum_{y \in X} w_y \langle \delta_x, P_v(\gamma) \delta_y \rangle \right|^2 \\
&= \sum_{x \in X} \left| \sum_{\gamma \in E} \sum_{y \in X} w_y \langle \delta_x, \lambda(\gamma) v \rangle \langle \lambda(\gamma) v, \delta_y \rangle \right|^2 \\
&\leq \sum_{x \in X} \left(\sum_{\gamma \in E} \sum_{y \in X} |w_y| \cdot |\langle \delta_x, \lambda(\gamma) v \rangle| \cdot |\langle \lambda(\gamma) v, \delta_y \rangle| \right)^2.
\end{aligned}$$

The scalar products in the last line are possibly non zero only if $\gamma^{-1}x, \gamma^{-1}y \in K$, so that $d(x, y) = d(\gamma^{-1}x, \gamma^{-1}y) \leq \text{diam } K =: d$, and $x, y \in \gamma K \subset EK := \cup_{\gamma' \in E} \gamma' K$. Let us observe that $|\{\gamma \in E : \gamma^{-1}x \in K\}| = |\{\gamma \in E : x \in \gamma K\}| \leq |K|$, so that

$$\begin{aligned}
\left\| \sum_{\gamma \in E} P_v(\gamma) w \right\|^2 &\leq \sum_{x \in X} \left(|\{\gamma \in E : \gamma^{-1}x \in K\}| \sum_{y \in B(x, d) \cap EK} |w_y| \right)^2 \\
&\leq |K|^2 \sum_{x \in X} \left(\sum_{y \in B(x, d) \cap EK} |w_y| \right)^2 \\
&\leq |K|^2 \sum_{x \in X} |B(x, d) \cap EK| \sum_{y \in B(x, d) \cap EK} |w_y|^2 \\
&\leq |K|^2 \sup_{x \in X} |B(x, d)| \sum_{y \in EK} \sum_{x \in B(y, d)} |w_y|^2 \\
&\leq |K|^2 (\sup_{x \in X} |B(x, d)|)^2 \sum_{y \in EK} |w_y|^2.
\end{aligned}$$

Let now H_ε be a finite subset of VX such that $\sum_{y \in X \setminus H_\varepsilon} |w_y|^2 < \varepsilon$, and set $F_\varepsilon := \{\gamma \in \Gamma : \gamma K \cap H_\varepsilon \neq \emptyset\}$, so that, for any finite set $E \subset \Gamma \setminus F_\varepsilon$, we have $EK = \cup_{\gamma \in E} \gamma K \subset \Gamma \setminus H_\varepsilon$, and $\sum_{y \in EK} |w_y|^2 < \varepsilon$, so that

$$\left\| \sum_{\gamma \in E} P_v(\gamma) w \right\|^2 \leq |K|^2 (\sup_{x \in X} |B(x, d)|)^2 \varepsilon,$$

which establishes the claim. \square

Theorem 5.2. *Let A be the adjacency matrix of a periodic graph X , π the trace representation of the von Neumann algebra $\mathcal{N}(X, \Gamma)$. Then*

$$\|A\| = \sup \sigma(\pi(A)) = \|\pi(A)\|,$$

i.e. $H := \|A\| - A$ does not have hidden spectrum.

Proof. Let v_n be the (normalised) Perron-Frobenius vector for the restriction A_n of A to the graph Λ_n , and let $P_{v_n}, P_{v_n}(\gamma)$ be as defined above. Moreover, for any

$m > n$, let us denote by E_m the projection on $\ell^2(\Lambda_m)$. Then

$$\begin{aligned}
\text{Tr}(E_m A P_{v_n}) &= \sum_{\gamma \in \Gamma} \text{Tr}(E_m A P_{v_n}(\gamma)) \\
&= \sum_{\gamma \in \Gamma} \langle \lambda(\gamma) v_n, E_m A \lambda(\gamma) v_n \rangle \\
&\geq \sum_{\gamma \in \Gamma} \langle \lambda(\gamma) v_n, E_m \lambda(\gamma) E_n \lambda(\gamma)^* A \lambda(\gamma) v_n \rangle \\
&= \|A_n\| \sum_{\gamma \in \Gamma} \langle \lambda(\gamma) v_n, E_m \lambda(\gamma) v_n \rangle \\
&= \|A_n\| \sum_{\gamma \in \Gamma} \text{Tr}(E_m P_{v_n}(\gamma)) = \|A_n\| \text{Tr}(E_m P_{v_n}),
\end{aligned}$$

where the inequality follows by the positivity of all the entries, and the last but one equality follows from the fact that $E_n \lambda(\gamma)^* A \lambda(\gamma) E_n = A_n$. As a consequence,

$$\|A_n\| \leq \frac{\text{Tr}(E_m A P_{v_n})}{\text{Tr}(E_m P_{v_n})} \rightarrow \frac{\tau(A P_{v_n})}{\tau(P_{v_n})} \leq \sup \sigma(\pi(A)).$$

where, for the last inequality, we used the following: let ξ_τ be the GNS vector, and $\eta = \frac{\pi(P_{v_n})\xi_\tau}{\|\pi(P_{v_n})\xi_\tau\|}$, then η is normalised, and $\frac{\tau(A P_{v_n})}{\tau(P_{v_n})} = \langle \eta, \pi(A)\eta \rangle$. The thesis follows. \square

Proposition 5.3. *Let X be a periodic graph. Then $H := \|A\| - A$ has finite critical density $\iff A$ is transient.*

Proof. The critical density of H is finite $\iff \tau(H^{-1}) = \int_0^\infty \frac{dN_H(\lambda)}{\lambda} < \infty$. Observe that $\tau(H^{-1}) = \sum_{x \in F} \langle \delta_x, H^{-1} \delta_x \rangle$, and recall that $\langle \delta_x, H^{-1} \delta_x \rangle < \infty$ for some $x \in X \iff \langle \delta_x, H^{-1} \delta_x \rangle < \infty$ for all $x \in X \iff A$ is transient. \square

6. THE SECULAR EQUATION

Our aim here is to show that additive perturbations of an essentially periodic graph can lead to hidden spectrum for the adjacency matrix of the perturbed graph. As explained above this is relevant for the occurrence of BEC condensation for the pure hopping model. As a starting point for this analysis we write down an eigenvalue equation, called secular equation, for the adjacency matrix in terms of objects associated to the unperturbed graph.

We start by considering very general perturbations of the graph under consideration, then we specialize the matter to the case of interest for our purposes. Let X, G be bounded degree graphs. We suppose we are adding and/or removing links from the graph X . Suppose further that G is (possibly) attached to X , describing another perturbation of the latter. Let Y be the resulting graph. Its adjacency matrix A_p can be written as

$$(6.1) \quad A_p = \begin{pmatrix} A + D & C \\ C^t & B \end{pmatrix},$$

where A, B are the adjacency matrices of X, G respectively, D describes the contribution of the links added and/or removed from X , and finally C describes the edges linking G to X . In the following $R_T(\lambda) := (\lambda I - T)^{-1}$ will denote the resolvent of the operator T defined for $\lambda \notin \sigma(T)$. From now on we suppose that the resulting

graph Y is of bounded degree, this implies that C and D are bounded operators. Define for $\lambda \notin \sigma(A) \cup \sigma(B)$,

$$(6.2) \quad S(\lambda) := (DR_A(\lambda) + CR_B(\lambda)C^t R_A(\lambda)) \upharpoonright_{\overline{\mathcal{R}(C) + \mathcal{R}(D)}}.$$

In the case under consideration, $S(\lambda) = PS(\lambda)P$ where P is the orthogonal projection on the closed subspace generated by the ranges of C and D , on which we suppose S naturally acts.

Theorem 6.1. *With the above notation, $\lambda \notin \sigma(A) \cup \sigma(B)$ is an eigenvalue of A_p iff 1 is an eigenvalue of $S(\lambda)$. If this is the case, the corresponding eigenvectors $v := \begin{pmatrix} x \\ y \end{pmatrix}$, respectively z , are related by*

$$(6.3) \quad \begin{aligned} z &= Dx + Cy, \\ x &= R_A(\lambda)z, \quad y = R_B(\lambda)C^t R_A(\lambda)z. \end{aligned}$$

Proof. Let $\lambda \notin \sigma(A) \cup \sigma(B)$, and suppose there is $v \in \ell^2(Y)$ such that $A_p v = \lambda v$. By (6.1), we recover from the first equation

$$(6.4) \quad x = R_A(\lambda)(Dx + Cy),$$

and, by multiplying both sides by D ,

$$(6.5) \quad Dx = DR_A(\lambda)(Dx + Cy).$$

Analogously, from the second equation, we obtain

$$(6.6) \quad y = R_B(\lambda)C^t x,$$

and from (6.4), by multiplying both sides by C ,

$$(6.7) \quad Cy = CR_B(\lambda)C^t x = CR_B(\lambda)C^t R_A(\lambda)(Dx + Cy).$$

Summing up (6.5) and (6.7), we obtain that $z = Dx + Cy$ is an eigenvector of $S(\lambda)$ corresponding to the eigenvalue 1. Conversely, suppose that z is an eigenvector of $S(\lambda)$ with eigenvalue 1, and $\lambda \notin \sigma(A) \cup \sigma(B)$. Define $v := \begin{pmatrix} R_A(\lambda)z \\ R_B(\lambda)C^t R_A(\lambda)z \end{pmatrix}$. Then, it is easy to show that v is an eigenvector of A_p with eigenvalue λ . \square

The equation

$$(6.8) \quad [DR_A(\lambda) + CR_B(\lambda)C^t R_A(\lambda)]z = z.$$

is called the *secular equation* in the present paper. It allows to compute the Perron–Frobenius eigenvalue of A_p in many cases of interest, including some infinite, density zero, additive perturbations of periodic graphs.⁴

Now we specialize the matter to the case of interest. Namely, we treat the case of finite additive perturbations of a (essentially) periodic graph. In this case, B, C, D are finite rank operators, with D positivity preserving, and acting on a finite dimensional subspace of $\ell^2(X)$. Thus, $S(\lambda)$ is a finite dimensional matrix whenever it is defined.

We observe that in principle, $\|A_p\|$ might not be an eigenvalue of A_p , even if it is always the maximum of $\sigma(A_p)$ (cf. the existence of generalized Perron–Frobenius eigenvectors, see [16]). However, if $\|A_p\| > \max\{\|A\|, \|B\|\}$, next result shows that $\|A_p\|$ is indeed an eigenvalue of A_p .

⁴Compare the computations in Section 9 with those in [4] used to prove the existence of the hidden spectrum for the comb graph.

Corollary 6.2. *Let Y be a finite perturbation of X , and A, B, C, D as above. If $\|A_p\| > \max\{\|A\|, \|B\|\}$ then 1 is an eigenvalue of $S(\|A_p\|)$ and $\|A_p\| \in \sigma_p(A_p)$.*

Proof. Let Y_n be an exhaustion of Y such that $Y_1 \supset G \cup \{x : \delta_x \in \mathcal{R}(C) + \mathcal{R}(D)\}$, $\|A_{Y_1}\| > \max\{\|A\|, \|B\|\}$, and set $X_n = Y_n \cap X$. By the results of Section 4 we may also assume that the Perron-Frobenius vectors $v_n = \begin{pmatrix} x_n \\ y_n \end{pmatrix}$ for A_{Y_n} , normalized by taking value 1 on a fixed vertex of Y_1 , converge pointwise to a generalized Perron-Frobenius vector $v = \begin{pmatrix} x \\ y \end{pmatrix}$ for A_Y . Applying Theorem 6.1 to any inclusion $X_n \subset Y_n$, we obtain that $z_n = Dx_n + Cy_n$ is an eigenvector with eigenvalue 1 for the matrix $S_n(\lambda_n) := (DR_{A_{X_n}}(\lambda_n) + CR_B(\lambda_n)C^t R_{A_{X_n}}(\lambda_n)) \upharpoonright_{\mathcal{R}(C) + \mathcal{R}(D)}$, with $\lambda_n = \|A_{Y_n}\|$. Let us observe that the vectors z_n belong to the same finite-dimensional vector space, on which all the matrices $S_n(\lambda_n)$ act. By construction, $\lim_n z_n = z := Dx + Cy$ and $\lim_n S_n(\lambda_n) = S(\|A_Y\|)$, namely z is an eigenvector with eigenvalue 1 of the matrix $S(\|A_Y\|)$. Applying again Theorem 6.1 we show that v is a true eigenvector of A_Y with eigenvalue $\|A_Y\|$. \square

We describe two particular cases of (6.8) when $D = 0$,⁵

$$(6.9) \quad R_B(\lambda)C^t R_A(\lambda)Cy = y, \quad x = R_A(\lambda)Cx;$$

and when $G = \emptyset$,

$$(6.10) \quad DR_A(\lambda)z = z, \quad x = R_A(\lambda)z.$$

Corollary 6.3. *Let Y be a finite perturbation of X , and A, B, C, D as above. Assume that $\|B\| \geq \|A\|$, or $\|B\| < \|A\|$ and A is recurrent. Then $\|A_p\| > \max\{\|A\|, \|B\|\}$ and $\|A_p\|$ is an eigenvalue of A_p .*

Proof. The function $\lambda \in (\max\{\|A\|, \|B\|\}, +\infty) \mapsto \|S(\lambda)\|$ is decreasing and tends to 0 when $\lambda \rightarrow +\infty$. When $\|B\| \geq \|A\|$, $\lim_{\lambda \rightarrow \|B\|} \|S(\lambda)\| = +\infty$. When $\|B\| < \|A\|$,

$\lim_{\lambda \rightarrow \|A\|} \|S(\lambda)\| = +\infty$ if and only if A is recurrent. So, in both cases, there exists a unique $\lambda_0 > \max\{\|A\|, \|B\|\}$ such that $\|S(\lambda_0)\| = 1$. Since $S(\lambda_0)$ has positive entries, 1 is an eigenvalue, whose eigenvector z , the Perron-Frobenius eigenvector, has positive entries. Applying Theorem 6.1 we get an eigenvector v of A_p , for the eigenvalue λ_0 , having positive entries. This immediately implies $\|A_p\| \geq \lambda_0 > \max\{\|A\|, \|B\|\}$, therefore, by the previous Corollary, $\|A_p\|$ is an eigenvalue, whose eigenvector v' has positive entries. Then $\lambda_0 = \|A_p\|$ and $v = v'$. \square

We end the present section by presenting a formula, which is needed in the sequel, which describes $R_{A_p}(\lambda)$ in terms of the resolvents R_A and R_B .

Proposition 6.4. *Let Y be a finite perturbation of X , and $A, B, C, D, S(\lambda)$ as above. Consider $\lambda \in \mathbb{C}$ such that $|\lambda| > \|A_p\|$, and choose $v = \begin{pmatrix} x \\ y \end{pmatrix}$ in $\ell^2(Y)$. Then*

$$(6.11) \quad R_{A_p}(\lambda)v = \begin{pmatrix} R_A(\lambda)(x + z) \\ R_B(\lambda)(C^t R_A(\lambda)x + y + C^t R_A(\lambda)z) \end{pmatrix},$$

where $z = (I - S(\lambda))^{-1}((DR_A(\lambda) + CR_B(\lambda)C^t R_A(\lambda))x + CR_B(\lambda)y)$.

If $|\lambda|$ is sufficiently large, the formula holds also for infinite, additive perturbations with density zero.

⁵Notice that the matrix $C^t R_A(\lambda)C$ is nonnull if the graph Y is supposed to be connected.

Proof. If we show that $I - S(\lambda)$ is invertible, the result follows from a straightforward calculation. If the perturbation is finite, $1 \notin \sigma_p(S(\lambda))$, otherwise, by Theorem 6.1, λ would belong to $\sigma_p(A_p)$, against the hypothesis $|\lambda| > \|A_p\|$. Since $S(\lambda)$ is a finite dimensional matrix, this means $1 \notin \sigma(S(\lambda))$, i.e. $I - S(\lambda)$ is invertible. In the case of infinite perturbation, observe that $S(\lambda) \rightarrow 0$ when $|\lambda| \rightarrow \infty$, and the thesis follows. \square

Remark 6.5. Assume that the perturbation consists only of some extra edges, without adding vertices; then $C = B = 0$, and the result above becomes

$$R_{A_p}(\lambda) = R_A(\lambda) + R_A(\lambda)(I - DR_A(\lambda)|_{\mathcal{R}(D)})^{-1}DR_A(\lambda).$$

In order to have the result for infinite perturbations it suffices that $|\lambda| > \|A\| + \|D\|$.

7. PERTURBATIONS OF PERIODIC GRAPHS

The present section is devoted to some general results involving density zero perturbations of (essentially) periodic graphs.

Proposition 7.1. *Let X be an essentially periodic graph, and Y a density zero perturbation of X . Suppose that $\|A_Y\| > \|A_X\|$. Then Y has hidden spectrum. In addition,*

$$(7.1) \quad \rho_c^Y(\beta) = \rho^X(\beta, \mu)$$

where $\mu = \|A_X\| - \|A_Y\| < 0$.

Proof. Let $g_n(x)$ be a continuous mollifier equal to 1 if $x \leq \varepsilon$ and 0 if $x \geq \varepsilon + 1/n$. We have, with $H := \|A_Y\| - A_Y$ and $\varepsilon < \|A_Y\| - \|A_X\|$,

$$N_H([0, \varepsilon]) = \lim_n \tau(g_n(\|A_Y\| - A_Y)) = \lim_n \tau(g_n(\|A_Y\| - A_X)) = 0.$$

As for the critical density, we have

$$\begin{aligned} \rho_c^Y(\beta) &\equiv \tau\left((e^{\beta(\|A_Y\|I - A_Y)} - 1)^{-1}\right) = \tau\left((e^{\beta(\|A_Y\|I - A_X)} - 1)^{-1}\right) \\ &= \tau\left((e^{\beta(\|A_X\|I - A_X - \mu I)} - 1)^{-1}\right) \equiv \rho^X(\beta, \mu). \end{aligned}$$

\square

Notice that (7.1) allows us to compute the critical density of the perturbed network by using the formula for the density of the unperturbed one. It is very interesting for physical applications, to compare such a BEC critical density (equivalently critical temperature) with the critical density (temperature) of the formation of the Baarden–Cooper pairs in the Josephson junctions.

We now consider finite subtractive perturbations of essentially periodic graphs.

Theorem 7.2. *Let Y be the graph obtained by removing a finite number of vertices and links from an essentially periodic graph X which does not have hidden spectrum. Then*

- (i) Y does not have hidden spectrum,
- (ii) the critical densities of X and Y are equal.

Proof. (i) Since X does not have hidden spectrum, and finite perturbations do not change the trace τ , it suffices to show that $\|A_Y\| = \|A_X\|$. It is known that $\|A_Y\| \leq \|A_X\|$.

We obtain

$$\begin{aligned} 0 &\leq \|A_X\| - \|A_Y\| \leq \|A_X\| - \|A_Y\| + E_m(\|A_Y\| - A_Y) \\ &= E_m(\|A_X\| - A_Y) = E_m(\|A_X\| - A_X) = 0. \end{aligned}$$

Therefore $\|A_Y\| = \|A_X\|$ and $E_m(\|A_Y\| - A_Y) = 0 = E_0(\|A_Y\| - A_Y)$, which is the claim.

(ii) Since $N_Y = N_X$, we obtain $\rho_Y(\beta, \mu) = \rho_X(\beta, \mu)$, and $\rho_c^Y(\beta) = \rho_c^X(\beta)$. \square

Notice that Theorem 7.2 holds true for zero density subtractive perturbations. In addition, it tells us that zero density subtractive perturbations do not alter the character of an essentially periodic graph, provided the graph under consideration does not exhibit hidden spectrum (e.g. a periodic graph). Theorem 7.2 generalizes a result in [10].

We now use the results of the previous section to show that very small additive perturbations of essentially periodic graphs provide examples of pure hopping low dimensional models with finite critical density.

Proposition 7.3. *Let X be an essentially-periodic graph with infinite critical density. Then, there exists a point $x_0 \in X$ such that if we add to X only one vertex $\#$ linked to x_0 , then the graph $X \cup \{\#\}$ has finite critical density.*

Proof. Denote by A the adjacency matrix of X . The secular equation (6.9) for $X \cup \{\#\}$ becomes $\lambda^{-1} \langle \delta_{x_0}, R_A(\lambda) \delta_{x_0} \rangle = 1$, x_0 being the vertex (to be determined) of X to which $\#$ is connected. We will show that there exists $\lambda > \|A\|$ satisfying the previous equation.

$$\begin{aligned} \infty &= \int_0^{\|A\|} \frac{dN_H(\lambda)}{\lambda} = \int_0^\infty \frac{d\tau(E_H(\lambda))}{\lambda} = \int_0^\infty \frac{d\tau(E_A(\|A\| - \lambda))}{\lambda} = \int_{-\infty}^{\|A\|} \frac{d\tau(E_A(\nu))}{\|A\| - \nu} \iff \\ &\int_0^{\|A\|} \frac{d\tau(E_A(\lambda))}{\|A\| - \lambda} = \infty. \end{aligned}$$

Therefore,

$$\begin{aligned} \infty &= \int_0^{+\infty} \frac{dN_A(a)}{\|A\| - a} = \int_0^{+\infty} \left(\sum_{n=0}^{+\infty} \|A\|^{-(n+1)} a^n \right) dN_A(a) \\ &= \sum_{n=0}^{+\infty} \|A\|^{-(n+1)} \int_0^{+\infty} a^n dN_A(a) \equiv \sum_{n=0}^{+\infty} \|A\|^{-(n+1)} \tau(A^n) \\ &\equiv \sum_{n=0}^{+\infty} \|A\|^{-(n+1)} \lim_{\Lambda_k \uparrow X} \left(\frac{1}{|\Lambda_k|} \sum_{x \in \Lambda_k} \langle \delta_x, A^n \delta_x \rangle \right) \\ &\leq \liminf_{\Lambda_k \uparrow X} \left(\frac{1}{|\Lambda_k|} \sum_{x \in \Lambda_k} \left(\sum_{n=0}^{+\infty} \|A\|^{-(n+1)} \langle \delta_x, A^n \delta_x \rangle \right) \right). \end{aligned}$$

Here, the second equality follows by the monotone convergence Theorem, and the last inequality by the Fatou Lemma.

Then there exists $x_0 \in X$ such that

$$\langle \delta_0, R_A(\|A\|) \delta_0 \rangle = \sum_{n=0}^{+\infty} \|A\|^{-(n+1)} \langle \delta_{x_0}, A^n \delta_{x_0} \rangle > \|A\|.$$

This means that, for the decreasing function $\lambda^{-1}\langle\delta_{x_0}, R_A(\lambda)\delta_{x_0}\rangle$,

$$\begin{aligned}\lim_{\lambda \rightarrow +\infty} \lambda^{-1}\langle\delta_{x_0}, R_A(\lambda)\delta_{x_0}\rangle &= 0, \\ \lim_{\lambda \rightarrow \|A\|} \lambda^{-1}\langle\delta_{x_0}, R_A(\lambda)\delta_{x_0}\rangle &> 1.\end{aligned}$$

Namely, there exists a (unique) $\lambda > \|A\|$ such that the secular equation for $X \cup \{\#\}$ is satisfied, or by Theorem 6.1, there is a (unique) $\lambda > \|A\|$ which is an eigenvalue of A_p , which implies $\|A_p\| \geq \lambda > \|A\|$. Therefore, $E_m(\|A_p\| - A_p) = E_m(\|A_p\| - A) = \|A_p\| - \|A\| + E_m(\|A\| - A) \geq \|A_p\| - \|A\| > 0$, that is $X \cup \{\#\}$ exhibits low energy hidden spectrum. \square

Theorem 7.4. *Let X be a essentially-periodic graph. Then, there exists a point $x_0 \in X$ such that if we add to X only one vertex $\#$ linked to x_0 , then the graph $Y := X \cup \{\#\}$ verifies $\|A_Y\| > \|A_X\|$. This implies Y has hidden spectrum, hence finite critical density.*

Proof. Assume the critical density of X is finite. If $\|A_X\| = \|A_{X \cup \{\#\}}\|$, then the critical densities of X and $X \cup \{\#\}$ are equal, so that $X \cup \{\#\}$ has finite critical density. If $\|A_X\| < \|A_{X \cup \{\#\}}\|$, then $X \cup \{\#\}$ has hidden spectrum and finite critical density.

If X has infinite critical density the result follows from the Proposition above. \square

Remark 7.5.

(i) The proof of Proposition 7.3 is based on the fact that there exists x_0 such that $\langle\delta_{x_0}, R_A(\|A\|)\delta_{x_0}\rangle$ is large enough. This is trivially true if X is periodic, since in that case infinite critical density is equivalent to recurrence, namely $\langle\delta_x, R_A(\|A\|)\delta_x\rangle = +\infty$ for any x .⁶

(ii) Notice that, the divergence of the same integral $\int \frac{dN_A(a)}{\|A\| - a}$, on the one hand is responsible of the infinite critical density for the unperturbed graph X , on the other hand allows us to conclude that graphs obtained by considering very small additive perturbations of X have finite critical density.

Up to now we have shown that small perturbations of a graph can produce hidden spectrum to the pure hopping Hamiltonian. The remarkable fact, pointed out in the following theorem, is that such a network cannot exhibit hidden spectrum if we choose as Hamiltonian of the model the Laplace operator Δ of the graph.

Theorem 7.6. *Let X be a periodic amenable graph, and Y a density zero perturbation of X . Let the Hamiltonian H be the Laplace operator on Y . Then H does not have hidden spectrum.*

Proof. Let Δ , and Δ_p be the Laplacian of X , and the perturbed graph Y , respectively. We have

$$0 \leq E_0(\Delta_p) \leq E_m(\Delta_p) \equiv E_m(\Delta) = 0,$$

where the last equality follows by [9]. \square

⁶See [1] for results related to the BEC and the computation of the "vacuum distribution" $d\mu(\lambda) = \langle\delta_0, dE_A(\lambda)\delta_0\rangle$ in some cases of interest in quantum probability, such as the comb graph.

8. ONE DIMENSIONAL EXAMPLES

In this section, we exhibit some examples of graphs which have hidden spectrum. To prove that, we use proposition 7.1, so we have to compare the norms of the adjacency operators of a graph and its perturbation. To compute the norms, we use Perron-Frobenius theory, and in particular the secular equation (6.8). We start by considering the eigenvalue equation for the adjacency operator on a linear chain. As we shall see, this gives rise to a difference equation whose solutions form a 2-dimensional space. Therefore two more data, such as the value on boundary points, determine the solution for the given eigenvalue, and a further datum determines the eigenvalue. In this way we can calculate the Perron-Frobenius eigenvector and eigenvalue on one-sided or two-sided linear chains with perturbations. Also, we can compute the matrix elements of the resolvent $R_A(\lambda)$ for one-sided or two-sided linear chains. Indeed, the vector $v = R_A(\lambda)\delta_x$ satisfies the equation $(\lambda - A)v = \delta_x$, namely the eigenvector equation with a perturbation. We compute some examples below.

Example 8.1 (Modified chain graphs). Suppose that the perturbation is on the left of a linear chain, and it determines the first components, denoted by (α_0, β_0) , of an eigenvector corresponding to an eigenvalue $\lambda > 2$ for the adjacency matrix. Denote the other components on the right as $(\alpha_1, \beta_1, \alpha_2, \beta_2, \dots)$, see figure 2. The remaining components on the right are the solution of the finite-difference system

$$(8.1) \quad \begin{pmatrix} \alpha_{n+1} \\ \beta_{n+1} \end{pmatrix} = \begin{pmatrix} -1 & \lambda \\ -\lambda & \lambda^2 - 1 \end{pmatrix} \begin{pmatrix} \alpha_n \\ \beta_n \end{pmatrix}.$$

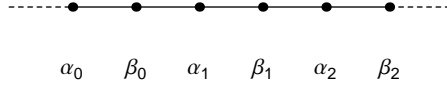


FIGURE 2. Infinite chain.

The eigenvalues of the matrix in (8.1) are

$$(8.2) \quad \mu_{\pm}(\lambda) = \frac{\lambda^2 - 2 \pm \lambda\sqrt{\lambda^2 - 4}}{2},$$

with corresponding eigenvectors

$$(8.3) \quad v_{\pm}(\lambda) := \begin{pmatrix} 2 \\ \lambda \pm \sqrt{\lambda^2 - 4} \end{pmatrix}.$$

Now we apply the previous considerations to the graph X in figure 3. Since the square-summable Perron-Frobenius eigenvector is unique, it is necessarily symmetric. So, we only search for symmetric eigenvectors of A_X . We have, in the previous notation,

$$(8.4) \quad \lambda = \alpha_0, \quad \lambda\alpha_0 = 1 + 2\beta_0.$$

By taking into account that, in order to get a square summable vector on the chain, $\begin{pmatrix} \alpha_0 \\ \beta_0 \end{pmatrix}$ cannot have a component along the eigenvector $v_+(\lambda)$, we obtain

$$(8.5) \quad \begin{pmatrix} \alpha_0 \\ \beta_0 \end{pmatrix} = a \begin{pmatrix} 2 \\ \lambda - \sqrt{\lambda^2 - 4} \end{pmatrix}.$$

Solving (8.4) and (8.5) w.r.t. λ , we obtain $\lambda = \sqrt{2 + \sqrt{5}} > 2$. Since the eigenvector we found has only positive components, it is the Perron-Frobenius eigenvector of A_X . Therefore $\|A_X\| = \lambda > 2 = \|A_{\mathbb{Z}}\|$, and, by proposition 7.1, the graph X has hidden spectrum.

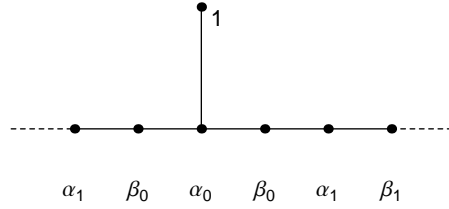


FIGURE 3. Infinite chain with a nail.

Other examples can be solved along the same lines, as those in figure 4. They both have hidden spectrum.



FIGURE 4. Some infinite graphs.

Another application of the previous method is to compute some matrix elements which are used in the sequel.

Proposition 8.2. *We have for the following matrix elements,*

$$(8.6) \quad \langle \delta_0, R_{A_{\mathbb{N}}}(\lambda) \delta_0 \rangle = \frac{2}{\lambda + \sqrt{\lambda^2 - 4}},$$

$$(8.7) \quad \langle \delta_0, R_{A_{box}}(\lambda) \delta_0 \rangle = \frac{2}{\lambda + \sqrt{\lambda^2 - 8}}$$

where A_{box} is the adjacency matrix of the box graph in figure 5.

Proof. We compute the latter, the computation of the former being similar. The one dimensional dynamical system (as that given in (8.1)) associated to the box graph is described by the matrix

$$\begin{pmatrix} -1 & \lambda \\ -\frac{\lambda}{2} & \frac{\lambda^2}{2} - 1 \end{pmatrix}$$

whose eigenvector associated to the negative eigenvalue has the form

$$\begin{pmatrix} 4 \\ \lambda - \sqrt{\lambda^2 - 8} \end{pmatrix}.$$

We have

$$\begin{pmatrix} \alpha_0 \\ \beta_0 \end{pmatrix} = a \begin{pmatrix} 4 \\ \lambda - \sqrt{\lambda^2 - 8} \end{pmatrix}, \quad \lambda\alpha_0 - 2\beta_0 = 1.$$

Solving w.r.t $\alpha_0 \equiv \langle \delta_0, R_{A_{box}}(\lambda)\delta_0 \rangle$, provides the assertion.

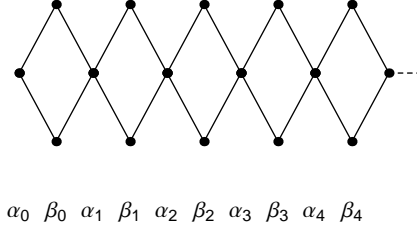


FIGURE 5. Box graph.

□

Remark 8.3. By applying the same calculation as before, we obtain

$$(8.8) \quad \langle \delta_0, R_{A_{\mathbb{Z}}}(\lambda)\delta_0 \rangle = \frac{1}{\sqrt{\lambda^2 - 4}}.$$

Now we apply the previous results to compute the Perron–Frobenius eigenvalue and/or eigenvector of some pivotal examples in order to show that they exhibit low dimensional hidden spectrum.

Remark 8.4. Even though \mathbb{N} is not a finite perturbation of a periodic graph, the disjoint union of two copies of \mathbb{N} , $\mathbb{N} \sqcup \mathbb{N}$, can be identified with the graph \mathbb{Z} with one link removed. The embedding of $\mathbb{N} \sqcup \mathbb{N}$ in \mathbb{Z} gives rise to an embedding of pairs (T, S) of operators in $\mathcal{A}_{FP}(\mathbb{N})$ into operators in $\mathcal{A}_{FP}(\mathbb{Z})$. We may therefore define the C*-algebra $\mathcal{A}(\mathbb{N})$ as that consisting of the operators $T \in \mathcal{A}_{FP}(\mathbb{N})$ such that $(T, T) \in \mathcal{A}(\mathbb{Z})$, endowed with the corresponding trace. This simple observation allows us to conclude that the critical density $\rho_c(\beta)$ of the graph \mathbb{N} , resp. the box-graph, is infinite as it coincides with that of \mathbb{Z} , resp. the bilateral box-graph (which is a \mathbb{Z} -lattice).

Example 8.5 (Star and star-box graphs). The star graph with $n \geq 3$ strands is composed by n copies of \mathbb{N} all connected to a single vertex $\#$, see figure 6, left. If $\lambda > 2 = \|A_{\mathbb{N}}\|$, the secular equation (6.9) is written as

$$(8.9) \quad \lambda = n \langle \delta_0, R_{A_{\mathbb{N}}}(\lambda)\delta_0 \rangle.$$

By taking into account (8.6), we obtain for the adjacency matrix A_n of the star graph with n strands,

$$\lambda \equiv \|A_n\| = \frac{n}{\sqrt{n-1}}.$$

Therefore, the star graph has hidden spectrum, see [4].

The star-box graph is made of $n \geq 3$ copies of the box graph connected to a single vertex $\#$, see figure 6, right. By taking into account (8.9) and (8.7), we have $\lambda = \frac{2n}{\lambda + \sqrt{\lambda^2 - 8}}$ which gives

$$\lambda = \frac{n}{\sqrt{n-2}}, \quad n \geq 4.$$

Since $\|A_{box}\| = 2\sqrt{2}$, the star-box graph has hidden spectrum iff $n \geq 5$.

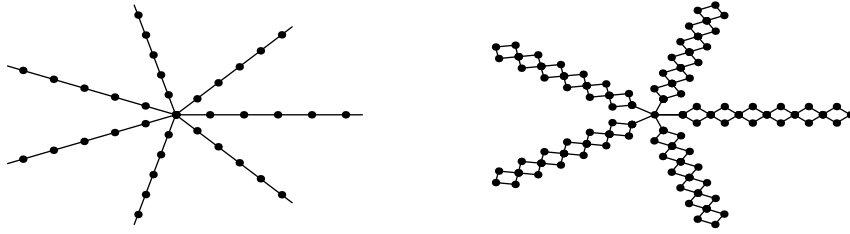


FIGURE 6. Star and star-box graphs.

Example 8.6 (Polygonal star and star-box graphs). We now apply the secular equation (6.10) to solve the star-box graph with n strands obtained by connecting $n \geq 3$ copies of the box graph through a polygon, see figure 7, right. The vector z in (6.10) is supported on the vertices of the polygon having n edges. If $\lambda > \|A_{box}\| = 2\sqrt{2}$, the secular equation for such a z living on the polygon, becomes

$$(8.10) \quad \langle \delta_0, R_{A_{box}}(\lambda) \delta_0 \rangle A_{pol} z = z,$$

where A_{box} , A_{pol} are the adjacency matrices of the box graph and the polygon respectively. Rotational invariance for the graph under consideration implies z has equal components, hence $2\langle \delta_0, R_{A_{box}}(\lambda) \delta_0 \rangle = 1$, independently of the number of the edges of the polygon. By taking into account (8.7), the previous equation gives $\lambda = 3 > 2\sqrt{2} \equiv \|A_{box}\|$. Hence, the star-box graph has hidden spectrum.

Another simple example of the star graph is that made of n strands connected by a polygon, see figure 7, left. By taking into account (8.10) and (8.6), we have $\frac{4}{\lambda + \sqrt{\lambda^2 - 4}} = 1$. Namely, $\lambda = \frac{5}{2} > 2 \equiv \|A_{\mathbb{N}}\|$, so that the graph has hidden spectrum.

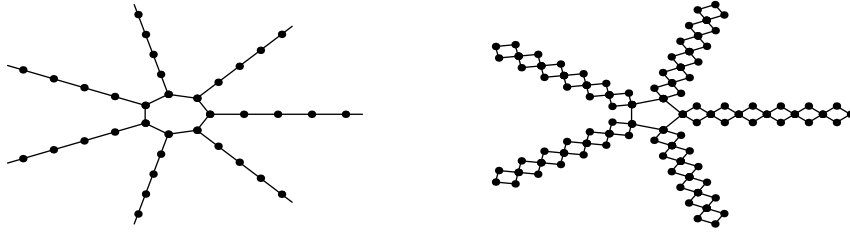


FIGURE 7. Polygonal star and star-box graphs.

Example 8.7 (H-graphs). We consider two copies of the bilateral infinite chain to which we add k links between the two origins, see figure 8. We call this graph an H -graph. We have

$$\alpha_0 \left(\frac{2}{\lambda - k} \right) = a\alpha_0 \left(\lambda + \sqrt{\lambda^2 - 4} \right) + b\alpha_0 \left(\lambda - \sqrt{\lambda^2 - 4} \right).$$

In order to have a square-summable eigenvector, we need $a = 0$, from which we obtain

$$\lambda = \sqrt{k^2 + 4}.$$

Therefore, the H -graph has hidden spectrum as soon as $k > 0$.

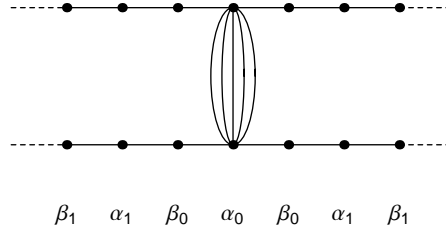


FIGURE 8. H graphs.

Example 8.8 (Modified ladder graphs). In the previous examples, we considered additive and subtractive perturbations separately. We now consider them together.

We consider the bilateral ladder graph modified as follows. We add $k - 1$ links at the origin, and remove $2n$ links symmetrically, see figure 9. We look at a modified ladder graph as a graph containing a suitable H -graph as a subgraph.

Since this graph contains the H -graph as a subgraph, $\|A_X\| \geq \sqrt{k^2 + 4}$. Since the ladder graph (*i.e.* for $k = 1$ and $n = 0$) has $\|A_{ladder}\| = 3$, the modified ladder graph X has hidden spectrum for all $k \geq 3$, $n \geq 0$. Moreover, by theorem 7.2 it follows that X has no hidden spectrum for $k = 0$ and any $n \geq 0$, or for $k = 1$ and any $n \geq 1$. Finally, it is possible to prove that, for $k = 2$, X has hidden spectrum for $n = 0$, and no hidden spectrum for $n \geq 1$.

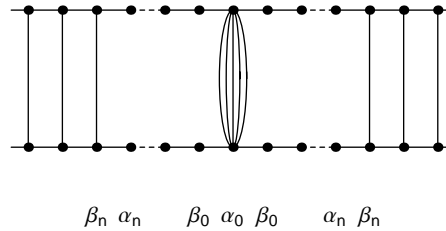


FIGURE 9. Modified ladder graphs.

9. COMB GRAPHS

In [4] the authors considered a graph, which they called comb graph, and showed that it has low energy hidden spectrum.

In general, one can define the comb product between two graphs, as in the following definition (cf. [1]).

Definition 9.1. Let G, H be graphs, and let $o \in VH$ be a given vertex. Then the comb product $X := G \dashv (H, o)$ is a graph with $VX := VG \times VH$, and $(g, h), (g', h') \in VX$ are adjacent *iff* $g = g'$ and $h \sim h'$ or $h = h' = o$ and $g \sim g'$. We call G the *base graph*, and H the *fibre graph*. When $o \in H$ is understood from the context, we omit it, and write $G \dashv H$.

In this and the next section we shall consider the so called comb graphs, with base \mathbb{Z}^d , $d \in \mathbb{N}$, and fibre \mathbb{Z} , with distinguished vertex $0 \in \mathbb{Z}$, and we denote them simply by $\mathbb{Z}^d \dashv \mathbb{Z}$. As we shall see, the comb graphs exhibit a different behaviour with respect to BEC, if $d \leq 2$ or $d \geq 3$.

The comb graph $\mathbb{Z} \dashv \mathbb{Z}$ can be described as an additive perturbation of the disconnected graph given by \mathbb{Z} copies of \mathbb{Z} (i.e. the fibres). The perturbation consists of adding some extra links: for any $n \in \mathbb{Z}$, there is a link connecting the zero-point of the n -th copy to the zero-point of the $(n+1)$ -th copy, see figure 1. The added links form a copy of \mathbb{Z} which is usually called the backbone in the Physics literature. We endow this graph with the regular exhaustion $\{\Lambda_n\}_{n \in \mathbb{N}}$, where Λ_n is the square $[-n, n] \times [-n, n]$. Here, and in the following, we denote by $[m, n] := \{z \in \mathbb{Z} : m \leq z \leq n\}$. In this sense the comb graph is a density zero perturbation of the disconnected graph given by infinitely many disjoint copies of \mathbb{Z} , indeed, according to Proposition 2.7, it is sufficient to note that

$$\frac{|(E(\mathbb{Z} \dashv \mathbb{Z}) \setminus E(\sqcup_{\mathbb{Z}} \mathbb{Z})) \cap E\Lambda_n|}{|V\Lambda_n|} = \frac{2n+1}{(2n+1)^2} \rightarrow 0.$$

As already said, it was shown in [4] that the comb graph $\mathbb{Z} \dashv \mathbb{Z}$ has low energy hidden spectrum. We generalize this result as follows.

Proposition 9.2. *The comb graph $G_d = \mathbb{Z}^d \dashv \mathbb{Z}$ has low energy hidden spectrum. In particular $\|A_{G_d}\| = 2\sqrt{d^2 + 1}$. The generalized Perron-Frobenius vector on G_d , obtained as point-wise limit of Perron-Frobenius vectors on Λ_n , can be explicitly calculated.*

Proof. Let us observe that Λ_n can be described as a finite comb graph $[-n, n]^d \dashv [-n, n]$, hence as a finite perturbation of the disjoint union of $(2n+1)^d$ copies of $[-n, n]$. Applying Theorem 6.1 with $\lambda_n = \|A_{[-n, n]^d \dashv [-n, n]}\|$, and in particular equation (6.10), we obtain the equation for the Perron-Frobenius eigenvector z_n for the matrix $S(\lambda_n)$ corresponding to the eigenvalue 1, given by

$$(9.1) \quad \langle \delta_0, R_{A_{[-n, n]}}(\lambda_n) \delta_0 \rangle A_{[-n, n]^d} z_n = z_n.$$

This means $R_{A_{[-n, n]}}(\lambda_n) \delta_0 \|A_{[-n, n]^d}\| = 1$. Taking the limit for $n \rightarrow \infty$, we obtain, with $\lambda := \lim \lambda_n$,

$$\langle \delta_0, R_{A_{\mathbb{Z}}}(\lambda) \delta_0 \rangle \|A_{\mathbb{Z}^d}\| = 1$$

which, according to (8.8), leads to

$$(9.2) \quad \frac{2d}{\sqrt{\lambda^2 - 4}} = 1,$$

or equivalently

$$\lambda = \|A_{G^d}\| = 2\sqrt{d^2 + 1}.$$

Notice that z_n is the Perron-Frobenius vector for $A_{[-n,n]^d}$, and we can normalize z_n in order to have value 1 in the origin. Then z_n converges pointwise to the unique generalized Perron-Frobenius vector for \mathbb{Z}^d , namely the vector which is constantly equal to 1 on \mathbb{Z}^d . Then, according to (6.10), the Perron-Frobenius vector x_n of $A_{[-n,n]^d \dashv [-n,n]}$ is given by $x_n = R_{A_{[-n,n]}}(\lambda_n)z_n$. As a consequence, for any $(\vec{j}, j) \in V(\mathbb{Z}^d \dashv \mathbb{Z})$,

$$\langle \delta_{\vec{j}, j}, x_n \rangle = \langle \delta_{\vec{j}, j}, R_{A_{[-n,n]}}(\lambda_n)z_n \rangle = \langle \delta_{\vec{j}}, z_n \rangle \langle \delta_j, R_{A_{[-n,n]}}(\lambda_n)\delta_0 \rangle.$$

Taking the limit for $n \rightarrow \infty$ we show that x_n converges pointwise; its limit is the generalized Perron-Frobenius vector x whose component (\vec{j}, j) is equal to $\langle \delta_j, R_{A_{\mathbb{Z}}}(\lambda)\delta_0 \rangle$, with $\lambda = 2\sqrt{d^2 + 1}$. \square

Let us consider the comb graph $G^d = \mathbb{Z}^d \dashv \mathbb{Z}$ together with the finite volume approximations $\Lambda_n = X_n \dashv Y_n$, where X_n is the graph $(\mathbb{Z}_{2n+1})^d$ (periodic boundary condition on the base graph), and Y_n is the finite chain $[-n, n]$.

Lemma 9.3. *Describing $\ell^2(\Lambda_n)$ as $\ell^2(X_n) \otimes \ell^2(Y_n)$, the Perron Frobenius eigenvector v_n for the adjacency operator A_{Λ_n} has the form $u_n \otimes R_{Y_n}(\|A_{\Lambda_n}\|)\delta_0$, where u_n is the vector constantly equal to 1 on X_n . Moreover,*

$$(9.3) \quad 2d\langle \delta_0, R_{Y_n}(\|A_{\Lambda_n}\|)\delta_0 \rangle = 1$$

Proof. Indeed, $A_{\Lambda_n} = I \otimes A_{Y_n} + A_{X_n} \otimes P_0$, P_0 denoting the one-dimensional projection on δ_0 , so that, for $v_n = u_n \otimes w_n$,

$$A_{\Lambda_n}v_n = u_n \otimes A_{Y_n}w_n + A_{X_n}u_n \otimes P_0w_n = u_n \otimes (A_{Y_n}w_n + 2dP_0w_n),$$

where we used the fact that the constant vector on X_n is the Perron-Frobenius vector for A_{X_n} , and the equality $\|A_{X_n}\| = 2d$.

Then, v_n is an eigenvector for A_{Λ_n} with eigenvalue t if $A_{Y_n}w_n + 2dP_0w_n - tw_n = 0$, which gives

$$(9.4) \quad (t - A_{Y_n})w_n = 2d\langle \delta_0, w_n \rangle \delta_0.$$

In particular, this implies

$$(9.5) \quad 2d\langle \delta_0, R_{Y_n}(t)\delta_0 \rangle = 1,$$

where $R_{Y_n}(t)$ denotes the resolvent $(t - A_{Y_n})^{-1}$. Let us note that the function $\langle \delta_0, R_{Y_n}(t)\delta_0 \rangle$ is decreasing in $(\|A_{Y_n}\|, +\infty)$, $\lim_{t \rightarrow \|A_{Y_n}\|} \langle \delta_0, R_{Y_n}(t)\delta_0 \rangle = +\infty$ and $\lim_{t \rightarrow +\infty} \langle \delta_0, R_{Y_n}(t)\delta_0 \rangle = 0$, therefore there exists a t for which condition (9.5) is satisfied. With such a t , the choice $w_n = R_{Y_n}(t)\delta_0$ gives rise to an eigenvector v_n by equation (9.4). Moreover such v_n has positive entries, hence is the Perron-Frobenius vector. This implies that the t satisfying equation (9.5) coincides with $\|A_{\Lambda_n}\|$, so it is unique. \square

We may now use the preceeding Lemma to obtain results on the graph G^d .

Lemma 9.4. *Let $v_n = u_n \otimes w_n$, with u_n constantly equal to 1 on the base X_n , $w_n = \|R_{Y_n}(\|A_{\Lambda_n}\|)\delta_0\|^{-1}R_{Y_n}(\|A_{\Lambda_n}\|)\delta_0$, and $v = u \otimes w$, with u constantly equal to 1 on the base graph \mathbb{Z}^d and $w = \|R_{\mathbb{Z}}(\|A\|)\delta_0\|^{-1}R_{\mathbb{Z}}(\|A\|)\delta_0$. Then*

(i) w_n converges in norm to w ,

(ii) v is a generalized Perron-Frobenius vector for A ,

(iii) $\langle \delta_0, R_{Y_n}(\lambda) \delta_0 \rangle = \frac{\tanh(n+1)\vartheta}{\sqrt{\lambda^2-4}}$, where $2 \cosh \vartheta = \lambda$,

(iv) $\|R_{\mathbb{Z}}(\|A\|) \delta_0\|^2 = \frac{\sqrt{d^2+1}}{4d^3}$.

Proof. (i) Let us observe that $\|A_{\Lambda_n}\| \rightarrow \|A\| = 2\sqrt{d^2+1}$, with $A = A_{G^d}$, while $\|A_{Y_n}\| \rightarrow \|A_{\mathbb{Z}}\| = 2$, hence Proposition 3.10 implies that $R_{Y_n}(\|A_{\Lambda_n}\|)$ converges strongly to $R_{\mathbb{Z}}(\|A\|)$.

(ii) The proof is analogous to the result in Section 4.

(iii) Setting $\lambda = 2 \cosh \vartheta$, it is not difficult to check that the vector $z(\lambda, n)$ defined by

(9.6)

$$z(\lambda, n)_j = \frac{\sinh[(n+1-|j|)\vartheta]}{2 \sinh \vartheta \cosh[(n+1)\vartheta]} = \tanh[(n+1)\vartheta] \frac{\cosh(|j|\vartheta)}{2 \sinh \vartheta} - \frac{\sinh(|j|\vartheta)}{2 \sinh \vartheta}, \quad |j| \leq n,$$

satisfies $(\lambda I - A_{Y_n})z(\lambda, n) = \delta_0$, therefore $\langle \delta_0, R_{Y_n}(\lambda) \delta_0 \rangle = z(\lambda, n)_0$. The thesis follows since $2 \sinh \vartheta = 2\sqrt{\cosh^2 \vartheta - 1} = \sqrt{\lambda^2 - 4}$.

(iv) By (i), $z(\|A\|, n)$ converges in norm to $R_{\mathbb{Z}}(\|A\|) \delta_0$, therefore

$$\langle \delta_j, R_{\mathbb{Z}}(\|A\|) \delta_0 \rangle = \frac{e^{-|j|\vartheta}}{2 \sinh \vartheta},$$

with $2 \cosh \vartheta = \|A\| = 2\sqrt{d^2+1}$. The thesis follows by a straightforward computation. \square

Let us set $v = \sum_{\vec{j}} \delta_{\vec{j}} \otimes v_{\vec{j}}$, $v_{\vec{j}}$ denoting the restriction of v to the fibre at the point \vec{j} . We say that $v \in \mathcal{S}_0$ if the sequence $\|v_{\vec{j}}\|$, $\vec{j} \in \mathbb{Z}^d$, is rapidly decreasing. We now show that $T_t \equiv e^{itH}$ defines a one-parameter group of Bogoliubov automorphisms on \mathcal{S}_0 .

Proposition 9.5. *Let $H = \|A\| - A$ on the comb graphs G_d . Then $e^{itH} \mathcal{S}_0 \subset \mathcal{S}_0$.*

Proof. If $v \in \ell^2(G_d)$ then

$$e^{itH} v = \frac{1}{2\pi i} \oint_{\gamma} e^{it(\|A\|-\lambda)} R_{G^d}(\lambda) v d\lambda$$

where γ is a Jordan curve surrounding counterclockwise the spectrum of H , and $R_{G^d}(\lambda) = (\lambda - A)^{-1}$ is the resolvent of the adjacency operator on the comb graph. Let $(\vec{j}, j) = (j_1, \dots, j_d, i)$ denote the coordinates of the comb graph G_d , and denote by $\delta_{\vec{j}} \otimes \delta_j$ the delta function on a point (\vec{j}, j) . Let us recall that, by Proposition 6.4, for $|\lambda|$ large enough, $R_{G^d}(\lambda) = I \otimes R_{\mathbb{Z}}(\lambda) + \Phi(\lambda) \otimes R_{\mathbb{Z}}(\lambda) P_0 R_{\mathbb{Z}}(\lambda)$ with

$$(9.7) \quad \Phi(\lambda) := s(\lambda) R_{\mathbb{Z}^d}(s(\lambda)) A_{\mathbb{Z}^d},$$

and $s(\lambda)$ is the holomorphic extension of $\sqrt{\lambda^2 - 4}$ to $\mathbb{C} \setminus [-2, 2]$.

Let us set $v = \sum_{\vec{j}} \delta_{\vec{j}} \otimes v_{\vec{j}}$, $v_{\vec{j}}$ denoting the restriction of v to the fibre at the point \vec{j} . Then the assumption $v \in \mathcal{S}_0$ amounts to say that the sequence $\|v_{\vec{j}}\|$, $\vec{j} \in \mathbb{Z}^d$, is rapidly decreasing, while the thesis, namely $e^{itH} v \in \mathcal{S}_0$, is equivalent to say that the sequence $\sum_{j \in \mathbb{Z}} |\langle \delta_{\vec{j}} \otimes \delta_j, e^{itH} v \rangle|^2$, $\vec{j} \in \mathbb{Z}^d$, is rapidly decreasing. By the equations above, we have

$$\begin{aligned} \langle \delta_{\vec{j}} \otimes \delta_j, R_{G^d}(\lambda) v \rangle &= \\ &= \langle \delta_j, R_{\mathbb{Z}}(\lambda) v_{\vec{j}} \rangle + \sum_{\vec{k} \in \mathbb{Z}^d} \langle \delta_{\vec{j}}, \Phi(\lambda) \delta_{\vec{k}} \rangle \langle \delta_j, R_{\mathbb{Z}}(\lambda) P_0 R_{\mathbb{Z}}(\lambda) v_{\vec{k}} \rangle. \end{aligned}$$

Choose the curve γ as a circle with radius greater than $\|A_{\mathbb{Z}}\| + \|A_{\mathbb{Z}^d}\| = 2(d+1)$, which surrounds the spectrum of A_{G_d} , and the hypothesis of Proposition 6.4 are satisfied, see Remark 6.5. With such a choice,

$$\begin{aligned} \sum_{j \in \mathbb{Z}} |\langle \delta_{\vec{j}} \otimes \delta_j, e^{itH} v \rangle|^2 &\leq \frac{\ell(\gamma)^2}{2\pi^2} \sup_{\lambda \in \gamma} |e^{-it\lambda}|^2 \|R_{\mathbb{Z}}(\lambda)\|^2 \|v_{\vec{j}}\|^2 + \\ &+ \frac{\ell(\gamma)^2}{2\pi^2} \sup_{\lambda \in \gamma} |e^{-it\lambda}|^2 \sum_{j \in \mathbb{Z}} \left| \sum_{\vec{k} \in \mathbb{Z}^d} \langle \delta_{\vec{j}}, \Phi(\lambda) \delta_{\vec{k}} \rangle \langle \delta_j, R_{\mathbb{Z}}(\lambda) P_0 R_{\mathbb{Z}}(\lambda) v_{\vec{k}} \rangle \right|^2. \end{aligned}$$

Clearly the first summand is rapidly decreasing in \vec{j} . Concerning the second summand, we have

$$\begin{aligned} &\sup_{\lambda \in \gamma} \sum_{j \in \mathbb{Z}} \left| \sum_{\vec{k} \in \mathbb{Z}^d} \langle \delta_{\vec{j}}, \Phi(\lambda) \delta_{\vec{k}} \rangle \langle \delta_j, R_{\mathbb{Z}}(\lambda) P_0 R_{\mathbb{Z}}(\lambda) v_{\vec{k}} \rangle \right|^2 = \\ &= \sup_{\lambda \in \gamma} \left\| \sum_{\vec{k} \in \mathbb{Z}^d} \langle \delta_{\vec{j}}, \Phi(\lambda) \delta_{\vec{k}} \rangle R_{\mathbb{Z}}(\lambda) P_0 R_{\mathbb{Z}}(\lambda) v_{\vec{k}} \right\|^2 \\ &\leq \sup_{\lambda \in \gamma} \|R_{\mathbb{Z}}(\lambda)\|^4 \left\| \sum_{\vec{k} \in \mathbb{Z}^d} \langle \delta_{\vec{j}}, \Phi(\lambda) \delta_{\vec{k}} \rangle v_{\vec{k}} \right\|^2. \end{aligned}$$

The thesis now amounts to show that, for any multi-index $\alpha = (\alpha_1, \dots, \alpha_d)$, the sequence $\sup_{\lambda \in \gamma} \|\vec{j}^\alpha \sum_{\vec{k} \in \mathbb{Z}^d} \langle \delta_{\vec{j}}, \Phi(\lambda) \delta_{\vec{k}} \rangle v_{\vec{k}}\|$ is bounded.

Passing to the Fourier transform on the torus \mathbb{T}^d , and setting $v(\vartheta) := \sum_{\vec{k} \in \mathbb{Z}^d} e^{i\vec{k}\vartheta} v_{\vec{k}}$, we have

$$\begin{aligned} &\sup_{\lambda \in \gamma} \left\| \vec{j}^\alpha \sum_{\vec{k} \in \mathbb{Z}^d} \langle \delta_{\vec{j}}, \Phi(\lambda) \delta_{\vec{k}} \rangle v_{\vec{k}} \right\| = \\ &= \sup_{\lambda \in \gamma} \left\| \vec{j}^\alpha \int_{\mathbb{T}^d} \sum_{\vec{k} \in \mathbb{Z}^d} e^{i(\mathbf{k}-\mathbf{j})\vartheta} \frac{2s(\lambda) \sum_{l=1}^d \cos 2\pi\vartheta_l}{s(\lambda) - 2 \sum_{l=1}^d \cos 2\pi\vartheta_l} v_{\vec{k}} d\vartheta \right\| \\ &= \sup_{\lambda \in \gamma} \left\| (-i)^{|\alpha|} \int_{\mathbb{T}^d} e^{-i\vec{j}\vartheta} \frac{\partial^{|\alpha|}}{\partial \vartheta^\alpha} \left(\frac{2s(\lambda) \sum_{l=1}^d \cos 2\pi\vartheta_l}{s(\lambda) - 2 \sum_{l=1}^d \cos 2\pi\vartheta_l} v(\vartheta) \right) d\vartheta \right\| \\ &\leq \sup_{\lambda \in \gamma} \sup_{\vartheta \in \mathbb{T}^d} \left\| \frac{\partial^{|\alpha|}}{\partial \vartheta^\alpha} \left(\frac{2s(\lambda) \sum_{l=1}^d \cos 2\pi\vartheta_l}{s(\lambda) - 2 \sum_{l=1}^d \cos 2\pi\vartheta_l} v(\vartheta) \right) \right\|. \end{aligned}$$

Since $v_{\vec{k}}$ is rapidly decreasing, $v(\vartheta) \in \mathcal{C}^\infty(\mathbb{T}^d, \ell^2(\mathbb{Z}))$. The thesis follows. \square

We end the present section by pointing out the following fact. For $v \in \mathcal{S}_0$, define on the Weyl operators $\alpha_t(W(v)) := W(T_t v)$. We obtain a one-parameter group of $*$ -automorphisms $t \mapsto \alpha_t$ on the CCR algebra $\text{CCR}(\mathcal{S}_0)$. Namely, $(\text{CCR}(\mathcal{S}_0), \alpha)$ is the dynamical system which is of interest in our context.

10. THERMODYNAMICAL STATES FOR COMB GRAPHS IN THE CONDENSATION REGIME

10.1. General results. We will consider here the comb $G^d = \mathbb{Z}^d \dashv \mathbb{Z}$ with the finite volume approximations $\Lambda_n = X_n \dashv Y_n$, where $X_n = (\mathbb{Z}_{2n+1})^d$ (periodic boundary condition on the base graph), and Y_n is the line graph $\{-n, \dots, 0, \dots, n\}$. Our aim is to study the states ω_n relative to the Gibbs grand canonical ensemble on the finite volume approximations Λ_n with chemical potential μ_n , and the existence of the limit state ω on the comb graph G^d in the condensation regime, that is when $\mu_n \rightarrow 0$. The case $\mu_n \rightarrow \mu < 0$ of non condensation regime, presents no further technical difficulties, and is described in Theorem 3.11. To avoid technicalities, we suppose $\mu_n < 0$.

Let A_{Λ_n} be the adjacency matrix of the comb graph Λ_n . The matrices A_{Λ_n} can all be considered as operators acting on $\ell^2(VG^d)$, if we identify $X_n \dashv Y_n$ with $[-n, n]^d \dashv [-n, n] \subset \mathbb{Z}^d \dashv \mathbb{Z}$. We have to study the limit behaviour of

$$\omega_n(a^+(\xi)a(\eta)), \quad \xi, \eta \in \mathbb{S}_0.$$

Let us denote by $H_n = (\|A\| - \mu_n)I - A_{\Lambda_n}$ the Hamiltonian on Λ_n with chemical potential μ_n . We want to compute the limit

$$(10.1) \quad \lim_n \langle \eta, (e^{\beta H_n} - I)^{-1} \xi \rangle,$$

for suitable vectors η, ξ . We first write

$$(e^{\beta H_n} - I)^{-1} = ((e^{\beta H_n} - I)^{-1} - (\beta H_n)^{-1}) + (\beta H_n)^{-1}.$$

Then

Lemma 10.1. *$(e^{\beta H_n} - I)^{-1} - (\beta H_n)^{-1}$ converges to $(e^{\beta H} - I)^{-1} - (\beta H)^{-1}$ in the strong operator topology.*

Proof. Indeed the function $(e^{\beta \lambda} - I)^{-1} - (\beta \lambda)^{-1}$ is continuous on $[0, \infty)$, hence the result follows by Proposition 3.10. \square

We have therefore reduced the computation of (10.1) to the computation of

$$(10.2) \quad \lim_n \langle \eta, (\beta H_n)^{-1} \xi \rangle.$$

Set $\lambda_n := \|A\| - \mu_n$. Since $H_n^{-1} = R_{\Lambda_n}(\lambda_n)$ by definition, the identity of Proposition 6.4 implies

$$(10.3) \quad H_n^{-1} = I \otimes R_{Y_n}(\lambda_n) + \Phi_n \otimes R_{Y_n}(\lambda_n) P_0 R_{Y_n}(\lambda_n),$$

where $P_0 = |\delta_0\rangle\langle\delta_0|$ (in bra-ket notation), and $\Phi_n = (I - \langle\delta_0, R_{Y_n}(\lambda_n)\delta_0\rangle A_{X_n})^{-1} A_{X_n}$. Let us notice that, by Lemma 9.4 (iii), $\langle\delta_0, R_{Y_n}(\lambda_n)\delta_0\rangle \rightarrow (2d)^{-1}$, hence we set

$$(10.4) \quad \langle\delta_0, R_{Y_n}(\lambda_n)\delta_0\rangle = \frac{1}{2(d + \varepsilon_n)},$$

and we have $\varepsilon_n \rightarrow 0$. Setting $\eta = \sum_{\vec{j}} \delta_{\vec{j}} \otimes \eta_{\vec{j}}$, $\xi = \sum_{\vec{k}} \delta_{\vec{k}} \otimes \xi_{\vec{k}}$ we have

$$(10.5) \quad \langle \eta, H_n^{-1} \xi \rangle = \langle \eta, I \otimes R_{Y_n}(\lambda_n) \xi \rangle + \sum_{\vec{j}, \vec{k}} \langle \delta_{\vec{j}}, \Phi_n \delta_{\vec{k}} \rangle \langle \eta_{\vec{j}}, R_{Y_n}(\lambda_n) P_0 R_{Y_n}(\lambda_n) \xi_{\vec{k}} \rangle.$$

We first observe that, by Proposition 3.10,

$$(10.6) \quad \lim_n \langle \eta, I \otimes R_{Y_n}(\lambda_n) \xi \rangle = \langle \eta, I \otimes R_{\mathbb{Z}}(2\sqrt{d^2 + 1}) \xi \rangle$$

Let us now compute $\lim_n \langle \delta_{\vec{j}}, \Phi_n \delta_{\vec{k}} \rangle$. Making use of discrete Fourier transform we get

$$\langle \delta_{\vec{j}}, \Phi_n \delta_{\vec{k}} \rangle = \begin{cases} 2d(d + \varepsilon_n) \int_{\mathbb{T}^d} \frac{(\frac{1}{d} \sum_{i=1}^d \cos \vartheta_i) \cos((\vec{j} - \vec{k})\vec{\vartheta})}{\varepsilon_n + \sum_{i=1}^d (1 - \cos \vartheta_i)} dm_n(\vec{\vartheta}) & \vec{j}, \vec{k} \in [-n, n]^d \\ 0 & \vec{j}, \vec{k} \notin [-n, n]^d \end{cases}$$

where m_n is the normalized measure on \mathbb{T}^d given by

$$(10.7) \quad m_n = \frac{1}{(2n+1)^d} \sum_{\vec{j} \in [-n, n]^d} \delta_{\frac{2\pi}{2n+1} \vec{j}}.$$

We now write

$$(10.8) \quad \int_{\mathbb{T}^d} \frac{(\frac{1}{d} \sum_{i=1}^d \cos \vartheta_i) \cos((\vec{j} - \vec{k})\vec{\vartheta})}{\varepsilon_n + \sum_{i=1}^d (1 - \cos \vartheta_i)} dm_n(\vec{\vartheta}) = k_n + Q_n(\delta_{\vec{j}}, \delta_{\vec{k}}),$$

where

$$(10.9) \quad k_n = \int_{\mathbb{T}^d} \frac{1}{\varepsilon_n + \sum_{i=1}^d (1 - \cos \vartheta_i)} dm_n(\vec{\vartheta}),$$

$$(10.10) \quad Q_n(\delta_{\vec{j}}, \delta_{\vec{k}}) = \int_{\mathbb{T}^d} \frac{(\frac{1}{d} \sum_{i=1}^d \cos \vartheta_i) \cos((\vec{j} - \vec{k})\vec{\vartheta}) - 1}{\varepsilon_n + \sum_{i=1}^d (1 - \cos \vartheta_i)} dm_n(\vec{\vartheta}),$$

According to Lemma 10.1 and equations (10.6), (10.5), (10.8), we have proved the following.

Lemma 10.2.

$$(10.11) \quad \begin{aligned} \langle \eta, (e^{\beta H_n} - I)^{-1} \xi \rangle &= \langle \eta, ((e^{\beta H_n} - I)^{-1} - (\beta H_n)^{-1}) \xi \rangle + \frac{1}{\beta} \langle \eta, I \otimes R_{Y_n}(\lambda_n) \xi \rangle \end{aligned}$$

$$(10.12) \quad + \frac{2d(d + \varepsilon_n)}{\beta} \sum_{\vec{j}, \vec{k} \in [-n, n]^d} Q_n(\delta_{\vec{j}}, \delta_{\vec{k}}) \langle \eta_{\vec{j}}, R_{Y_n}(\lambda_n) P_0 R_{Y_n}(\lambda_n) \xi_{\vec{k}} \rangle$$

$$(10.13) \quad + \frac{2d(d + \varepsilon_n)}{\beta} k_n \sum_{\vec{j}, \vec{k} \in [-n, n]^d} \langle \eta_{\vec{j}}, R_{Y_n}(\lambda_n) P_0 R_{Y_n}(\lambda_n) \xi_{\vec{k}} \rangle,$$

We now show that $Q_n(\delta_{\vec{j}}, \delta_{\vec{k}})$ converges to

$$(10.14) \quad Q(\delta_{\vec{j}}, \delta_{\vec{k}}) := \int_{\mathbb{T}^d} \frac{(\frac{1}{d} \sum_{i=1}^d \cos \vartheta_i) \cos((\vec{j} - \vec{k})\vec{\vartheta}) - 1}{\sum_{i=1}^d (1 - \cos \vartheta_i)} dm(\vec{\vartheta}),$$

where dm denotes the normalized Lebesgue measure on \mathbb{T}^d .

Proposition 10.3.

$$|Q_n(\delta_{\vec{j}}, \delta_{\vec{k}}) - Q(\delta_{\vec{j}}, \delta_{\vec{k}})| \leq \alpha_n (1 + |\vec{j} - \vec{k}|^2)$$

for a suitable infinitesimal sequence α_n .

Proof. Indeed, let $g(\varepsilon, \vec{\nu}, \vec{\vartheta}) = \frac{(\frac{1}{d} \sum_{i=1}^d \cos \vartheta_i) \cos(\vec{\nu} \vec{\vartheta}) - 1}{\varepsilon + \sum_{i=1}^d (1 - \cos \vartheta_i)}$, with $\vec{\nu} := \vec{j} - \vec{k}$, and write

$$\begin{aligned} |Q_n(\delta_{\vec{j}}, \delta_{\vec{k}}) - Q(\delta_{\vec{j}}, \delta_{\vec{k}})| &= \left| \int_{\mathbb{T}^d} \left(g(\varepsilon_n, \vec{\nu}, \vec{\vartheta}) dm_n(\vec{\vartheta}) - g(0, \vec{\nu}, \vec{\vartheta}) dm(\vec{\vartheta}) \right) \right| \\ &\leq \left| \int_{B(0,r)} \left(g(\varepsilon_n, \vec{\nu}, \vec{\vartheta}) dm_n(\vec{\vartheta}) - g(0, \vec{\nu}, \vec{\vartheta}) dm(\vec{\vartheta}) \right) \right| \\ &\quad + \int_{B(0,r)^c} |g(\varepsilon_n, \vec{\nu}, \vec{\vartheta}) - g(0, \vec{\nu}, \vec{\vartheta})| dm_n(\vec{\vartheta}) \\ &\quad + \left| \int_{B(0,r)^c} g(0, \vec{\nu}, \vec{\vartheta}) \left(dm_n(\vec{\vartheta}) - dm(\vec{\vartheta}) \right) \right| \end{aligned}$$

We now fix $\delta > 0$, and observe that we may find r independent of $\varepsilon_n, \vec{\nu}$ and $\vec{\vartheta}$ such that the first summand of the r.h.s. above is bounded by $\delta(1 + |\vec{\nu}|^2)$. Moreover, on $\mathbb{T}^d \setminus B(0, r)$, $|g(\varepsilon_n, \vec{\nu}, \vec{\vartheta}) - g(0, \vec{\nu}, \vec{\vartheta})| \leq \frac{4\varepsilon_n}{r^4} < \delta$ for sufficiently big n . Finally $|\nabla g(0, \vec{\nu}, \vec{\vartheta})| \leq \frac{4}{r^4}(1 + |\vec{\nu}|)$ on $\mathbb{T} \setminus B(0, r)$, hence the third summand is bounded by $\delta(1 + |\vec{\nu}|)$ for sufficiently big n . The thesis follows. \square

In order to continue the analysis of the limit state ω , we have to study $\lim_n k_n$, where k_n was defined in (10.9). As we shall see, this requires to study the low-dimensional case and the high-dimensional case separately.

10.2. The comb graph with low-dimensional base graph.

Proposition 10.4. *If $d \leq 2$, then $k_n \rightarrow +\infty$.*

Proof.

$$k_n \geq \int_{\mathbb{T}^d} \frac{1}{\varepsilon_n + \frac{|\vec{\vartheta}|^2}{2}} dm_n(\vec{\vartheta}) \geq \int_{\mathbb{T}^d} \frac{1}{\varepsilon_n + \frac{|\vec{\vartheta}|^2}{2}} dm(\vec{\vartheta}) = \text{const} \int_0^\pi \frac{\vartheta^{d-1}}{\varepsilon_n + \frac{\vartheta^2}{2}} d\vartheta$$

The thesis follows since the last integral diverges for $d = 1, 2$ when $\varepsilon_n \rightarrow 0$. \square

Now we may prove the main result of this subsection.

Theorem 10.5. *Let $\xi, \eta \in \mathbb{S}_0$. For each sequence $\mu_n \rightarrow 0$,*

$$(10.15) \quad \omega_n(a^+(\xi)a(\eta)) = k'_n \langle \xi, v_n \rangle \langle v_n, \eta \rangle + C_n(\xi, \eta),$$

where $k'_n \rightarrow +\infty$, $\langle v_n, \eta \rangle \rightarrow \langle v, \eta \rangle$ for any $\eta \in \mathbb{S}_0$, with $v_n \in V\Lambda_n$ and $v_n \rightarrow v$, the generalized Perron–Frobenius vector on G^d described in Lemma 9.4, and C_n converges to a sesquilinear form with domain containing \mathbb{S}_0 .

Proof. Let us observe that, with $w_n := \|R_{Y_n}(\lambda_n)\delta_0\|^{-1} R_{Y_n}(\lambda_n)\delta_0$,

$$\begin{aligned} \langle \eta_{\vec{j}}, R_{Y_n}(\lambda_n) P_0 R_{Y_n}(\lambda_n) \xi_{\vec{k}} \rangle &= \langle \eta_{\vec{j}}, R_{Y_n}(\lambda_n) \delta_0 \rangle \langle R_{Y_n}(\lambda_n) \delta_0, \xi_{\vec{k}} \rangle \\ &= \|R_{Y_n}(\lambda_n) \delta_0\|^2 \langle \eta_{\vec{j}}, w_n \rangle \langle w_n, \xi_{\vec{k}} \rangle. \end{aligned}$$

Denote by $v_n := u_n \otimes w_n$, where, here and in the following, we use the definitions in Lemma 9.4 for the vectors u_n, u, w, v .

On the one hand, we get

$$\sum_{\vec{j}, \vec{k}} \langle \eta_{\vec{j}}, R_{Y_n}(\lambda_n) P_0 R_{Y_n}(\lambda_n) \xi_{\vec{k}} \rangle = \|R_{Y_n}(\lambda_n) \delta_0\|^2 \langle \eta, v_n \rangle \langle v_n, \xi \rangle,$$

and

$$\begin{aligned} \langle v - v_n, \xi \rangle &= \left| \sum_{\vec{j}} \langle u, \delta_{\vec{j}} \rangle \langle w, \xi_{\vec{j}} \rangle - \langle u_n, \delta_{\vec{j}} \rangle \langle w_n, \xi_{\vec{j}} \rangle \right| \\ &\leq \sum_{\vec{j} \notin [-n, n]^d} \|\xi_{\vec{j}}\| + \|w - w_n\| \sum_{\vec{j} \in [-n, n]^d} \|\xi_{\vec{j}}\| \end{aligned}$$

which tends to 0, since $w_n \rightarrow w$ in norm, and the sum $\sum_{\vec{j}} \|\xi_{\vec{j}}\|$ is finite, since ξ is in \mathcal{S}_0 . We have thus proved that the term (10.13) in Lemma 10.2 gives the first summand in the statement, with $k'_n = \beta^{-1} 2d(d + \varepsilon_n) k_n \|R_{Y_n}(\lambda_n) \delta_0\|^2$.

Again by Proposition 3.10, the term (10.11) in Lemma 10.2 converges to

$$\langle \eta, ((e^{\beta H} - I)^{-1} - (\beta H)^{-1}) \xi \rangle + \frac{1}{\beta} \langle \eta, I \otimes R_{\mathbb{Z}}(\lambda_\infty) \xi \rangle,$$

where $\lambda_\infty := 2\sqrt{d^2 + 1}$. As for the term (10.12) in Lemma 10.2, we want to show that it converges to

$$\frac{2d^2}{\beta} \sum_{\vec{j}, \vec{k}} Q(\delta_{\vec{j}}, \delta_{\vec{k}}) \langle \eta_{\vec{j}}, R_{\mathbb{Z}}(\lambda_\infty) P_0 R_{\mathbb{Z}}(\lambda_\infty) \xi_{\vec{k}} \rangle,$$

or, equivalently, that

$$\sum_{\vec{j}, \vec{k} \in [-n, n]^d} Q_n(\delta_{\vec{j}}, \delta_{\vec{k}}) \langle \eta_{\vec{j}}, w_n \rangle \langle w_n, \xi_{\vec{k}} \rangle \rightarrow \sum_{\vec{j}, \vec{k}} Q(\delta_{\vec{j}}, \delta_{\vec{k}}) \langle \eta_{\vec{j}}, w \rangle \langle w, \xi_{\vec{k}} \rangle$$

where Q was defined in (10.14). Indeed

$$\begin{aligned} &\left| \sum_{\vec{j}, \vec{k}} Q(\delta_{\vec{j}}, \delta_{\vec{k}}) \langle \eta_{\vec{j}}, w \rangle \langle w, \xi_{\vec{k}} \rangle - \sum_{\vec{j}, \vec{k} \in [-n, n]^d} Q_n(\delta_{\vec{j}}, \delta_{\vec{k}}) \langle \eta_{\vec{j}}, w_n \rangle \langle w_n, \xi_{\vec{k}} \rangle \right| \\ &\leq \sum_{\vec{j}, \vec{k} \in [-n, n]^d} |Q(\delta_{\vec{j}}, \delta_{\vec{k}}) - Q_n(\delta_{\vec{j}}, \delta_{\vec{k}})| \|\eta_{\vec{j}}\| \|\xi_{\vec{k}}\| \\ &+ \sum_{\vec{j}, \vec{k} \notin [-n, n]^d} |Q(\delta_{\vec{j}}, \delta_{\vec{k}})| \|\eta_{\vec{j}}\| \|\xi_{\vec{k}}\| \\ &+ \sum_{\vec{j}, \vec{k} \in [-n, n]^d} |Q_n(\delta_{\vec{j}}, \delta_{\vec{k}})| |\langle \eta_{\vec{j}}, w \rangle \langle w, \xi_{\vec{k}} \rangle - \langle \eta_{\vec{j}}, w_n \rangle \langle w_n, \xi_{\vec{k}} \rangle|. \end{aligned}$$

According to Proposition 10.3, the first summand on the r.h.s. is majorized by

$$\alpha_n \sum_{\vec{j}, \vec{k}} (1 + |\vec{j} - \vec{k}|)^2 \|\eta_{\vec{j}}\| \|\xi_{\vec{k}}\|$$

which tends to zero since α_n does and the sum is finite because $\eta, \xi \in \mathcal{S}_0$.

Since, by definition of Q , $|Q(\delta_{\vec{j}}, \delta_{\vec{k}})| \leq (1 + |\vec{j} - \vec{k}|)^2$, the second summand is majorized by

$$\sum_{\vec{j}, \vec{k} \notin [-n, n]^d} (1 + |\vec{j} - \vec{k}|)^2 \|\eta_{\vec{j}}\| \|\xi_{\vec{k}}\|$$

which tends to zero since the sum is finite as above.

In the last summand, we again have $|Q_n(\delta_{\vec{j}}, \delta_{\vec{k}})| \leq (1 + |\vec{j} - \vec{k}|)^2$, and

$$|\langle \eta_{\vec{j}}, w \rangle \langle w, \xi_{\vec{k}} \rangle - \langle \eta_{\vec{j}}, w_n \rangle \langle w_n, \xi_{\vec{k}} \rangle| \leq 2\|w - w_n\| \|\eta_{\vec{j}}\| \|\xi_{\vec{k}}\|.$$

Therefore the last summand is bounded by

$$2\|w - w_n\| \sum_{\vec{j}, \vec{k}} (1 + |\vec{j} - \vec{k}|)^2 \|\eta_{\vec{j}}\| \|\xi_{\vec{k}}\|,$$

hence tends to 0 as before. \square

Remark 10.6.

(i) Theorem 10.5 tells us that it is impossible to construct any locally normal states describing BEC (i.e. whenever $\mu_n \rightarrow 0$) on the combs $\mathbb{Z}^d \wr \mathbb{Z}$, $d = 1, 2$, see Lemma 3.2. In addition, the divergence of the two-point function depends only on the amount of condensate (cf. C_n in Theorem 10.5 converges to a sesquilinear form which is finite on \mathbb{S}_0).

(ii) Let us observe that the lack of a locally normal thermodynamic state describing condensation does not mean that physically the condensation does not occur. It indeed means that in nonhomogeneous networks, particles condensate even in the configuration space, due to the shape of the wave function of the ground state. Then more and more particles tend to lay in the low energy spectrum and along the base space. The system cannot accommodate them.

10.3. The comb graph with high-dimensional base graph. We will show here that, for the comb $G^d = \mathbb{Z}^d \wr \mathbb{Z}$, $d \geq 3$, it is possible to construct infinite volume locally normal KMS states. As before, we will consider the finite volume approximations $\Lambda_n = X_n \wr Y_n$ of G^d , where $X_n = (\mathbb{Z}_{2n+1})^d$ (periodic boundary condition on the base graph), and Y_n is the line graph $\{-n, \dots, 0, \dots, n\}$. Our aim is to show that, for a carefully chosen sequence of chemical potentials $\mu_n \rightarrow 0$, we obtain locally normal thermodynamical states exhibiting Bose-Einstein condensation.

Recalling the discussion above, the estimate from below of k_n in Proposition 10.4 does not imply that $k_n \rightarrow \infty$ when $d \geq 3$. In order to describe its behavior we have to split k_n in two parts, the zero component of the integral and the rest, $k_n = k_n^0 + k_n^+$, with

$$(10.16) \quad k_n^0 = \frac{1}{(2n+1)^d \varepsilon_n},$$

$$(10.17) \quad k_n^+ = \int_{\mathbb{T}^d \setminus \{\vec{0}\}} \frac{1}{\varepsilon_n + \sum_{i=1}^d (1 - \cos \vartheta_i)} dm_n(\vec{\vartheta}).$$

Our aim is to show that

$$\lim_n \int_{\mathbb{T}^d \setminus \{\vec{0}\}} \frac{1}{\varepsilon_n + \sum_{i=1}^d (1 - \cos \vartheta_i)} dm_n(\vec{\vartheta}) = \int_{\mathbb{T}^d} \frac{1}{\sum_{i=1}^d (1 - \cos \vartheta_i)} dm(\vec{\vartheta}).$$

Let us set $\varphi(\vec{\vartheta}) = \sum_{j=1}^d (1 - \cos \vartheta_j)$. The integrand is positive, hence, for any $\varepsilon > 0$ we have, for n large enough,

$$\begin{aligned} \int_{\mathbb{T}^d \setminus \{\vec{0}\}} \frac{dm_n(\vec{\vartheta})}{\varepsilon + \varphi(\vec{\vartheta})} &\leq \int_{\mathbb{T}^d \setminus \{\vec{0}\}} \frac{dm_n(\vec{\vartheta})}{\varepsilon_n + \varphi(\vec{\vartheta})} \leq \int_{\mathbb{T}^d \setminus \{\vec{0}\}} \frac{dm_n(\vec{\vartheta})}{\varphi(\vec{\vartheta})} \\ \downarrow n \rightarrow \infty \\ \int_{\mathbb{T}^d} \frac{dm(\vec{\vartheta})}{\varepsilon + \varphi(\vec{\vartheta})} &\xrightarrow{\varepsilon \rightarrow 0} \int_{\mathbb{T}^d} \frac{dm(\vec{\vartheta})}{\varphi(\vec{\vartheta})} \end{aligned}$$

The result will then follow if

$$\int_{\mathbb{T}^d \setminus \{\vec{0}\}} \frac{1}{\varphi(\vec{\vartheta})} (dm_n(\vec{\vartheta}) - dm(\vec{\vartheta})) \rightarrow 0.$$

Such integral can be rewritten as

$$\int_{\mathbb{T}^d \setminus \{\vec{0}\}} \left(\frac{1}{\varphi(\vec{\vartheta})} - \frac{1}{|\vec{\vartheta}|^2} \right) (dm_n(\vec{\vartheta}) - dm(\vec{\vartheta})) + \int_{\mathbb{T}^d \setminus \{\vec{0}\}} \frac{1}{|\vec{\vartheta}|^2} (dm_n(\vec{\vartheta}) - dm(\vec{\vartheta})).$$

The first term tends to zero because of the following Lemma.

Lemma 10.7. *Let F be a bounded function on $[-\pi, \pi]^d$ which is continuous but in zero. Then*

$$\int_{[-\pi, \pi]^d \setminus \{\vec{0}\}} F(\vec{\vartheta}) dm_n(\vec{\vartheta}) \rightarrow \int_{[-\pi, \pi]^d} F(\vec{\vartheta}) dm(\vec{\vartheta})$$

Proof. Let $\|F\|_\infty = C$, and, for a given $\varepsilon > 0$, choose δ such that $(2\delta)^d C < \varepsilon/2$. Clearly F is uniformly continuous on $([-\delta, \delta]^d)^c$, therefore there exists an n such that

$$\left| \int_{([--\delta, \delta]^d)^c} F(\vec{\vartheta}) dm_n(\vec{\vartheta}) - \int_{([--\delta, \delta]^d)^c} F(\vec{\vartheta}) dm(\vec{\vartheta}) \right| < \varepsilon/2.$$

With such choice

$$\left| \int_{[-\pi, \pi]^d \setminus \{\vec{0}\}} F(\vec{\vartheta}) dm_n(\vec{\vartheta}) - \int_{[-\pi, \pi]^d} F(\vec{\vartheta}) dm(\vec{\vartheta}) \right| < \varepsilon,$$

and the Lemma is proved. \square

The second term tends to 0 as shown by the following Lemma.

Lemma 10.8. *Let $d \geq 3$. Then*

$$\lim_n \int_{\mathbb{T}^d \setminus \{\vec{0}\}} \frac{1}{|\vec{\vartheta}|^2} dm_n(\vec{\vartheta}) = \int_{\mathbb{T}^d} \frac{1}{|\vec{\vartheta}|^2} dm(\vec{\vartheta}).$$

Proof. Since the domains of our integrals are contained in $[-\pi, \pi]^d$, we may assume

$$m_n = \frac{1}{(2n+1)^d} \sum_{\vec{j} \in \mathbb{Z}^d} \delta_{\frac{2\pi}{2n+1} \vec{j}}.$$

Denote by $C_{\vec{j}}$ the square

$$\prod_{i=1}^d \left[\frac{2\pi}{2n+1} \vec{j}_i, \frac{2\pi}{2n+1} (\vec{j}_i + 1) \right)$$

and by $\vec{1}$ the vector with components constantly equal to 1, then, for any \vec{j} such that $j_i > 0$, $i = 1, \dots, d$, we have

$$(10.18) \quad \int_{C_{\vec{j}+\vec{1}}} \frac{1}{|\vec{\vartheta}|^2} dm_n(\vec{\vartheta}) \leq \int_{C_{\vec{j}}} \frac{1}{|\vec{\vartheta}|^2} dm(\vec{\vartheta}) \leq \int_{C_{\vec{j}}} \frac{1}{|\vec{\vartheta}|^2} dm_n(\vec{\vartheta}).$$

As a consequence we have

$$\int_{\mathbb{T}^d} \frac{1}{|\vec{\vartheta}|^2} dm(\vec{\vartheta}) = 2^d \lim_{n \rightarrow \infty} \int_{[\frac{2\pi}{2n+1}, \pi]^d} \frac{1}{|\vec{\vartheta}|^2} dm(\vec{\vartheta}) \leq \lim_{n \rightarrow \infty} \int_{\mathbb{T}^d \setminus \{\vec{0}\}} \frac{1}{|\vec{\vartheta}|^2} dm_n(\vec{\vartheta}).$$

We now prove the opposite inequality. We decompose the lattice

$$\left\{ \frac{2\pi}{2n+1} \vec{j} : \vec{j} \in \mathbb{Z}^d \right\} \cap \mathbb{T}^d$$

according to the number of non-zero components of \vec{j} . Therefore, setting $\mathbb{T}_n^{d+} = \left\{ \frac{2\pi}{2n+1} \vec{k} \in (0, \pi]^d : \vec{k} \in \mathbb{Z}^d \right\}$, we get

$$(10.19) \quad \int_{\mathbb{T}^d \setminus \{\vec{0}\}} \frac{1}{|\vec{\vartheta}|^2} dm_n(\vec{\vartheta}) = \sum_{j=1}^d \binom{d}{j} 2^j (2n+1)^{j-d} \int_{\mathbb{T}_n^{j+}} \frac{1}{|\vec{\vartheta}|^2} dm_n(\vec{\vartheta}).$$

If $j \geq 3$, we use inequality (10.18) to get

$$(10.20) \quad (2n+1)^{j-d} \int_{\mathbb{T}_n^{j+}} \frac{1}{|\vec{\vartheta}|^2} dm_n(\vec{\vartheta}) \leq (2n+1)^{j-d} \int_{[0, \pi]^j} \frac{1}{|\vec{\vartheta}|^2} dm(\vec{\vartheta})$$

and note that the r.h.s. tends to zero if $j < d$.

For $j = 1, 2$, we decompose \mathbb{T}_n^{j+} as $\left\{ \frac{2\pi}{2n+1} \vec{1} \right\} \cup \left\{ \frac{2\pi}{2n+1} \vec{k} \in (0, \pi]^d : \vec{k} \neq \vec{1} \right\}$. Hence, again by inequality (10.18), we get

$$\begin{aligned} (2n+1)^{j-d} \int_{\mathbb{T}_n^{j+}} \frac{1}{|\vec{\vartheta}|^2} dm_n(\vec{\vartheta}) &\leq \frac{1}{(2n+1)^d} \left(\frac{2n+1}{2\pi} \right)^2 \frac{1}{j} \\ &\quad + 2^{-j} (2n+1)^{j-d} \int_{\mathbb{T}^j \setminus B(0, \frac{2\pi}{2n+1})} \frac{1}{|\vec{\vartheta}|^2} dm(\vec{\vartheta}). \end{aligned}$$

Both summands on the r.h.s. are infinitesimal. Therefore equations (10.19) and (10.20) give

$$\lim_n \int_{\mathbb{T}^d \setminus \{\vec{0}\}} \frac{1}{|\vec{\vartheta}|^2} dm_n(\vec{\vartheta}) \leq \int_{\mathbb{T}^d} \frac{1}{|\vec{\vartheta}|^2} dm(\vec{\vartheta})$$

□

We have proved the following Lemma.

Lemma 10.9.

$$\lim_n k_n^+ = \int_{\mathbb{T}^d} \frac{1}{\sum_{i=1}^d (1 - \cos \vartheta_i)} dm(\vec{\vartheta}).$$

This result, together with the definition of Φ_n and Proposition 10.3, gives

Proposition 10.10.

$$(10.21) \quad (2d(d + \varepsilon_n))^{-1} \langle \delta_{\vec{j}}, \Phi_n \delta_{\vec{k}} \rangle = \begin{cases} k_n^0 + k_n^+ + Q_n(\delta_{\vec{j}}, \delta_{\vec{k}}) & \vec{j}, \vec{k} \in [-n, n]^d, \\ 0 & \vec{j}, \vec{k} \notin [-n, n]^d. \end{cases}$$

There exists an infinitesimal sequence α_n'' such that

$$(10.22) \quad |(k_n^+ + Q_n(\delta_{\vec{j}}, \delta_{\vec{k}})) - (2d^2)^{-1} \langle \delta_{\vec{j}}, \Phi \delta_{\vec{k}} \rangle| \leq \alpha_n'' (1 + |\vec{j} - \vec{k}|^2),$$

where $\Phi := 2d R_{\mathbb{Z}^d}(2d)A_{\mathbb{Z}^d}$, so that, by Fourier transform,

$$(10.23) \quad \langle \delta_{\vec{j}}, \Phi \delta_{\vec{k}} \rangle = 2d^2 \int_{\mathbb{T}^d} \frac{(\frac{1}{d} \sum_{i=1}^d \cos \vartheta_i) \cos((\vec{j} - \vec{k}) \vec{\vartheta})}{\sum_{i=1}^d (1 - \cos \vartheta_i)} dm(\vec{\vartheta}).$$

10.4. The choice of μ_n . In order to have a finite limit for k_n^0 we have to assume that $\lim_n ((2n+1)^d \varepsilon_n)^{-1}$ is finite. The following holds.

Lemma 10.11.

$$\lim_n \frac{1}{-\mu_n \text{vol}(X_n)} = c < \infty \Leftrightarrow \lim_n k_n^0 = \lim_n \frac{1}{\varepsilon_n (2n+1)^d} = c \frac{2d}{\sqrt{d^2+1}} < \infty.$$

Proof. By (10.4) and Lemma 9.4 (iii), we obtain $\varepsilon_n = \frac{\sqrt{\lambda_n^2 - 4}}{2 \tanh(n+1)\tau_n} - d$, with $2 \cosh \tau_n = \lambda_n = \|A_{G^d}\| - \mu_n$, and $\|A_{G^d}\| = \lambda_\infty = 2\sqrt{d^2+1}$. Then

$$\frac{1}{\varepsilon_n (2n+1)^d} = \left(\frac{\sqrt{\lambda_n^2 - 4} - 2d}{2 \tanh(n+1)\tau_n} (2n+1)^d + d \frac{1 - \tanh(n+1)\tau_n}{\tanh(n+1)\tau_n} (2n+1)^d \right)^{-1}.$$

Since $\tau_n \rightarrow \tau_\infty$ with $\cosh \tau_\infty = \sqrt{d^2+1}$, $\tanh(n+1)\tau_n$ tends to 1 exponentially fast, and the second summand above tends to 0. As for the first summand,

$$\left(\frac{\sqrt{\lambda_n^2 - 4} - 2d}{2 \tanh(n+1)\tau_n} (2n+1)^d \right)^{-1} = \left(\frac{(-2\mu_n \lambda_\infty + \mu_n^2)(2n+1)^d}{2(\tanh(n+1)\tau_n)(\sqrt{\lambda_n^2 - 4} + 2d)} \right)^{-1}$$

and the latter has a finite limit if and only if $(-\mu_n(2n+1)^d)^{-1}$ has a finite limit. The thesis follows since $\text{vol}(X_n) = (2n+1)^d$. \square

Lemma 10.12. *Let us assume $\lim_n (-\mu_n \text{vol}(X_n))^{-1} = c < \infty$. Then, for any $\eta, \xi \in \mathcal{S}_0$,*

$$\lim_n \langle \eta, (H_n)^{-1} \xi \rangle = \langle \eta, I \otimes R_{\mathbb{Z}}(\lambda_\infty) \xi \rangle + \sum_{\vec{j}, \vec{k}} \langle \delta_{\vec{j}}, \Phi \delta_{\vec{k}} \rangle \langle \eta_{\vec{j}}, w \rangle \langle w, \xi_{\vec{k}} \rangle + c \langle \eta, v \rangle \langle v, \xi \rangle.$$

Proof. Let us recall that, by Lemma 10.2,

$$\begin{aligned} \langle \eta, H_n^{-1} \xi \rangle &= \langle \eta, I \otimes R_{Y_n}(\lambda_n) \xi \rangle \\ &\quad + 2d(d+\varepsilon) \sum_{\vec{j}, \vec{k} \in [-n, n]^d} (Q_n(\delta_{\vec{j}}, \delta_{\vec{k}}) + k_n^+) \langle \eta_{\vec{j}}, R_{Y_n}(\lambda_n) P_0 R_{Y_n}(\lambda_n) \xi_{\vec{k}} \rangle \\ &\quad + 2d(d+\varepsilon) k_n^0 \sum_{\vec{j}, \vec{k} \in [-n, n]^d} \langle \eta_{\vec{j}}, R_{Y_n}(\lambda_n) P_0 R_{Y_n}(\lambda_n) \xi_{\vec{k}} \rangle. \end{aligned}$$

The first summand tends to $\langle \eta, I \otimes R_{\mathbb{Z}}(\lambda_\infty) \xi \rangle$ by Proposition 3.10. The second summand tends to $\sum_{\vec{j}, \vec{k}} \langle \delta_{\vec{j}}, \Phi \delta_{\vec{k}} \rangle \langle \eta_{\vec{j}}, w \rangle \langle w, \xi_{\vec{k}} \rangle$ as in the proof of Theorem 10.5, with the aid of (10.22).

By Lemma 10.11 and by the proof of Theorem 10.5, the last summand tends to

$$c \frac{4d^3}{\sqrt{d^2+1}} \|R_{\mathbb{Z}}(\lambda_\infty) \delta_0\|^2 \langle \eta, v \rangle \langle v, \xi \rangle.$$

The thesis follows by Lemma 9.4 (iv). \square

We also have the following.

Lemma 10.13. *When G^d is considered as an infinite perturbation of infinitely many disjoint fibres \mathbb{Z} , equation (6.11) is valid for any $\lambda > \|A\|$.*

Proof. Indeed the identity becomes

$$R_A(\lambda) = I \otimes R_{\mathbb{Z}}(\lambda) + (I - \langle \delta_0, R_{\mathbb{Z}}(\lambda) \delta_0 \rangle A_{\mathbb{Z}^d})^{-1} A_{\mathbb{Z}^d} \otimes R_{\mathbb{Z}}(\lambda) P_0 R_{\mathbb{Z}}(\lambda),$$

where $P_0 = |\delta_0\rangle\langle\delta_0|$. Such identity has been proved to hold whenever the operator $(I - \langle \delta_0, R_{\mathbb{Z}}(\lambda) \delta_0 \rangle A_{\mathbb{Z}^d})$ is invertible.

Since $\langle \delta_0, R_{\mathbb{Z}}(\lambda) \delta_0 \rangle = (\lambda^2 - 4)^{-1/2}$, we have $(I - \langle \delta_0, R_{\mathbb{Z}}(\lambda) \delta_0 \rangle A_{\mathbb{Z}^d}) = (\lambda^2 - 4)^{-1/2}(\sqrt{\lambda^2 - 4} - A_{\mathbb{Z}^d})$, which is invertible whenever $\sqrt{\lambda^2 - 4} > \|A_{\mathbb{Z}^d}\| = 2d$, namely whenever $\lambda > 2\sqrt{d^2 + 1} = \|A\|$. \square

From the last Lemma, we have

$$\begin{aligned} \langle \xi, R_A(\lambda) \eta \rangle &= \sum_{\vec{j}} \langle \eta_{\vec{j}}, R_{\mathbb{Z}}(\lambda) \xi_{\vec{j}} \rangle \\ &\quad + \sum_{\vec{j}, \vec{k}} \langle \delta_{\vec{j}}, (I - \langle \delta_0, R_{\mathbb{Z}}(\lambda) \delta_0 \rangle A_{\mathbb{Z}^d})^{-1} A_{\mathbb{Z}^d} \delta_{\vec{k}} \rangle \langle \eta_{\vec{j}}, R_{\mathbb{Z}}(\lambda) P_0 R_{\mathbb{Z}}(\lambda) \xi_{\vec{k}} \rangle. \end{aligned}$$

Taking the limit for $\lambda \rightarrow \|A\|^+$ in the equation above, and using Lemma 10.12, we conclude

$$(10.24) \quad \lim_n \langle \eta, H_n^{-1} \xi \rangle = \langle \eta, H^{-1} \xi \rangle + c \langle \eta, v \rangle \langle v, \xi \rangle.$$

10.5. Conclusion.

Theorem 10.14. *Let $d \geq 3$, G^d denote the comb $\mathbb{Z}^d \uparrow \mathbb{Z}$, $\Lambda_n = X_n \uparrow Y_n$ be its approximation. Moreover, $H_n = (\|A\| - \mu_n)I - A_{\Lambda_n}$ denotes the Hamiltonian on Λ_n with chemical potential μ_n , $H = \|A\|I - A$ denotes the pure hopping Hamiltonian on G^d , and v denotes the Perron-Frobenius generalized vector for A considered in Lemma 9.4. Then:*

- (i) *the adjacency operator A for G^d is transient,*
- (ii) $\mathbb{S}_0 \subset \mathcal{D}(R_{G^d}(\|A\|^{1/2}))$.

Assume now

$$(10.25) \quad \lim_n \frac{1}{-\mu_n \text{vol}(X_n)} = c < +\infty.$$

Then

$$(10.26) \quad \lim_n \langle \eta, (e^{\beta H_n} - I)^{-1} \xi \rangle = \langle \eta, (e^{\beta H} - I)^{-1} \xi \rangle + \frac{c}{\beta} \langle \eta, v \rangle \langle v, \xi \rangle \quad \forall \xi, \eta \in \mathbb{S}_0,$$

$$(10.27) \quad \lim_n \tau_{\Lambda_n}((e^{\beta H_n} - I)^{-1}) = \tau((e^{\beta H} - I)^{-1}),$$

where τ_{Λ_n} is the normalized trace on Λ_n .

Proof. Statement (ii) means that, for any $\xi \in \mathbb{S}_0$, $\lim_{\lambda \rightarrow \|A\|^+} \langle \xi, R_{G^d}(\lambda) \xi \rangle$ is finite, and this follows from Lemmas 10.12, and 10.13. Clearly (ii) \Rightarrow (i) when $\xi = \delta_J$.

Equation (10.26) summarizes the results proved above in this section.

We now prove equation (10.27). First we use Proposition 3.5 to show that, as in Lemma 10.1, we only have to compute $\lim_n \tau_{\Lambda_n}(H_n^{-1})$. Using formula (10.3), we have

$$\tau_{\Lambda_n}(H_n^{-1}) = \tau_{[-n, n]}(R_{Y_n}(\lambda_n)) + \tau_{[-n, n]^d}(\Phi_n) \tau_{[-n, n]}(R_{Y_n}(\lambda_n) P_0 R_{Y_n}(\lambda_n))$$

where $\tau_{[-n,n]^j}$ denotes the normalized trace on $[-n,n]^j$. Again by Proposition 3.5 we get $\tau_{[-n,n]}(R_{Y_n}(\lambda_n)) \rightarrow \tau_{\mathbb{Z}}(R_{\mathbb{Z}}(\|A\|))$ and $\tau_{[-n,n]}(R_{Y_n}(\lambda_n)P_0R_{Y_n}(\lambda_n)) \rightarrow \tau_{\mathbb{Z}}(R_{\mathbb{Z}}(\|A\|)P_0R_{\mathbb{Z}}(\|A\|)) = 0$, since $R_{\mathbb{Z}}(\|A\|)P_0R_{\mathbb{Z}}(\|A\|)$ is finite rank. The result follows if we show that $\tau_{[-n,n]^d}(\Phi_n)$ is bounded. Let us recall that, by (10.21), $\langle \delta_{\vec{j}}, \Phi_n \delta_{\vec{k}} \rangle = 2d(d + \varepsilon_n)(k_n^0 + k_n^+ + Q_n(\delta_{\vec{j}}, \delta_{\vec{k}}))$. By formula (10.22),

$$|\tau_{[-n,n]^d}(k_n^+ + Q_n)| \leq \alpha_n'' + (2d^2)^{-1}\tau_{[-n,n]^d}(\Phi),$$

while k_n^0 is bounded by (10.25). \square

Remark 10.15.

(i) According to the Theorem above, in order to get a finite contribution for the condensate in the two-point function, condition $|\mu_n| \geq \text{const } n^{-d}$ should be satisfied, for a suitable positive constant (cf. Lemma 10.11). In this case, again by the previous theorem, the condensate does not contribute to the density. This is because the condensate is spatially distributed according to the Perron-Frobenius vector, namely around the base graph, therefore the condensate in Λ_n grows as n^d , while the volume grows as n^{d+1} .

(ii) Conversely, if we try to construct the thermodynamical state as a limit with fixed density, in particular choosing the inverse temperature $\beta > 0$, a parameter $k > 0$, and μ_n in such a way that

$$(10.28) \quad \rho_{\Lambda_n}(\beta, \mu_n) = \rho_c(\beta) + k,$$

we do not get a finite two-point function on local vectors.

Indeed, according to the proof of equation (10.27), the only term there which depends on the sequence μ_n is $\tau_{[-n,n]^d}(\Phi_n P_{u_n})$, or, equivalently, $\tau_{\Lambda_n}(H_n^{-1} P_{v_n})$. We have

$$\tau_{\Lambda_n}(H_n^{-1} P_{v_n}) = (\|A\| - \mu_n - \|A_{\Lambda_n}\|)^{-1}(2n+1)^{-d-1},$$

which, together with equation (10.28), gives $\lim_n |\mu_n|^{-1}(2n+1)^{-d-1} = k$. But with this choice $\langle \eta, H_n^{-1} P_{v_n} \xi \rangle$ behaves like $|\mu_n|^{-1}(2n+1)^{-d} \langle \eta, v_n \rangle \langle v_n, \xi \rangle$, which diverges as soon as $\langle \eta, v_n \rangle \langle v_n, \xi \rangle \neq 0$.

(iii) A non locally normal infinite-volume KMS state can be always constructed on the space $\mathcal{D}_0 := \bigcup_{\varepsilon > 0} P_{\varepsilon} \ell^2(G_d)$, P_{ε} being the spectral projection of the Hamiltonian $H = \|A\| - A$ corresponding to the spectral subspace $[\varepsilon, +\infty)$. The two-point function is given for $X, Y \in \mathcal{D}_0$,

$$\omega(a^+(\xi)a(\eta)) = \langle \xi, (e^{\beta(\|A\| - A)} - 1)^{-1} \eta \rangle.$$

This can be obtained as infinite volume limit of any sequence of finite volume Gibbs states based on any sequence of chemical potentials $\mu_n \rightarrow 0$. As the elements of \mathcal{D}_0 are formally orthogonal to the Perron Frobenius eigenvector, no amount of condensate can be appreciated in such non locally normal state.

We end the present section by showing that the locally normal states described in the previous theorem are KMS for the dynamics generated on $\text{CCR}(\mathcal{S}_0)$ by the one-parameter group of Bogoliubov transformations e^{itH} .

Theorem 10.16. *If $d \geq 3$ then the states $\omega^{(c)}$, with two-point function*

$$\omega^{(c)}(a^+(\xi)a(\eta)) = \langle \eta, (e^{\beta H} - I)^{-1} \xi \rangle + \frac{c}{\beta} \langle \eta, v \rangle \langle v, \xi \rangle,$$

on the CCR algebra $\text{CCR}(\mathcal{S}_0)$ associated to \mathcal{S}_0 are β -KMS w.r.t. the time evolution α_t induced on \mathfrak{A} by $T_t := e^{itH}$.

Proof. By Proposition 9.5, $T_t \mathcal{S}_0 = \mathcal{S}_0$, thus it induces a one-parameter group of automorphisms of \mathfrak{A} , by putting

$$\alpha_t(W(\xi)) := W(T_t \xi).$$

By taking into account the form of the two-point function for the $\omega^{(c)}$ (cf. [3], pag. 79), they are automatically KMS, provided that the functions

$$t \in \mathbb{R} \mapsto \omega^{(c)}(W(\xi)\alpha_t(W(\eta))), \quad \xi, \eta \in \mathcal{S}_0$$

are continuous. Notice that (cf. [3], pag. 42)

$$(10.29) \quad \omega^{(c)}(W(\xi)) = \exp \left\{ -\frac{\|\xi\|^2}{4} \right\} \exp \left\{ -\frac{\omega^{(c)}(a^+(\xi)a(\xi))}{2} \right\},$$

and, by using the commutation rule,

$$(10.30) \quad \omega^{(c)}(W(\xi)\alpha_t(W(\eta))) = e^{-i\text{Im}\langle \xi, T_t \eta \rangle} \omega^{(c)}(W(\xi + T_t \eta)).$$

Thus, by taking into account (10.29), (10.30), it is enough to show that

$$t \in \mathbb{R} \mapsto \omega^{(c)}(a^+(\xi + T_t \eta)a(\xi + T_t \eta))$$

is continuous whenever $\xi, \eta \in \mathcal{S}_0$. Indeed,

$$\begin{aligned} \omega^{(c)}(a^+(\xi + T_t \eta)a(\xi + T_t \eta)) &= \\ &= \langle \xi + \eta, v \rangle \langle v, \xi + \eta \rangle + \langle \xi + T_t \eta, (e^{\beta H} - 1)^{-1}(\xi + T_t \eta) \rangle \\ &= |\langle v, \xi + \eta \rangle|^2 + \langle \xi, (e^{\beta H} - 1)^{-1} \xi \rangle + \langle T_t(e^{\beta H} - 1)^{-1/2} \eta, (e^{\beta H} - 1)^{-1/2} \xi \rangle + \\ &\quad + \langle (e^{\beta H} - 1)^{-1/2} \xi, T_t(e^{\beta H} - 1)^{-1/2} \eta \rangle + \langle \eta, (e^{\beta H} - 1)^{-1} \eta \rangle \end{aligned}$$

The proof follows as $\xi \in \mathcal{S}_0$ implies that $\xi \in \mathcal{D}((e^{\beta H} - 1)^{-1/2})$, as shown in Theorem 10.14 (ii). \square

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DIPARTIMENTO DI MATEMATICA, UNIVERSITÀ DI ROMA TOR VERGATA, VIA DELLA RICERCA SCIENTIFICA 1, ROMA 00133, ITALY

E-mail address: fidaleo@mat.uniroma2.it, guido@mat.uniroma2.it, isola@mat.uniroma2.it